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## Motivations

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- ➔ Because, apart from emitted elementary particles, **they carry the only information** that the experimental instruments can measure.
- ❖ Making clusters is **not an easy task**, because it involves, in a complex environment:
  - ▶ the fundamental nuclear properties,
  - ▶ quantum effects,
  - ▶ and variable timescales.



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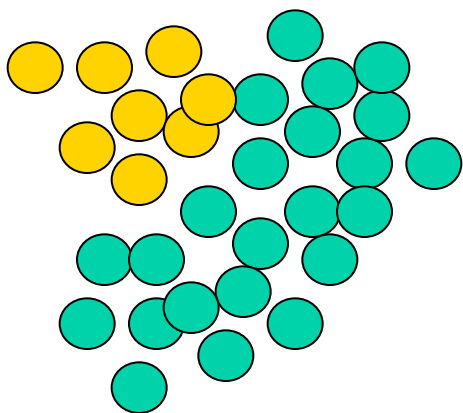
Simulations show: Clusters chosen that way at early times are the pre-fragments of the final state clusters, because fragments are not a random collection of nucleons at the end but initial-final state correlations.

# SACA: How does this work?

Simulated Annealing Procedure: PLB301:328,1993; later called SACA.

## 2 steps:

1) Pre-select good «candidates» for fragments according to proximity criteria: real space coalescence = Minimum Spanning Tree (MST) procedure.



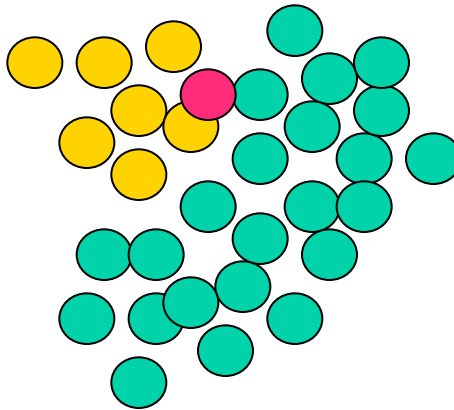
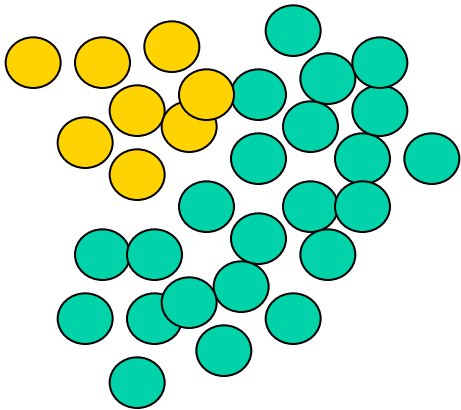
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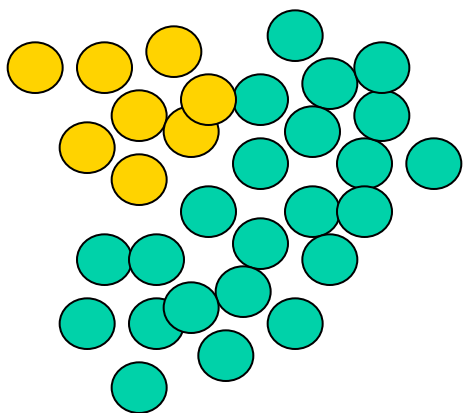
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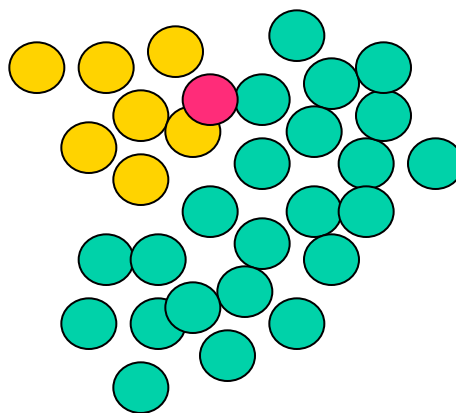
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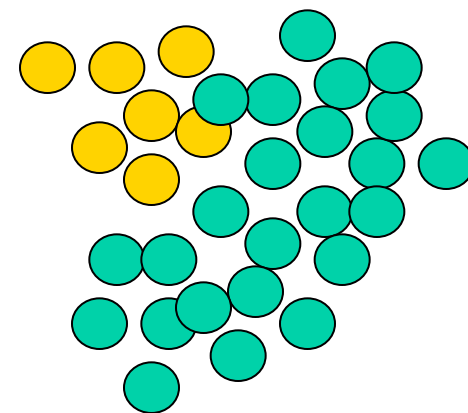
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$$E = E_{kin}^1 + E_{kin}^2 + V^1 + V^2$$



$$E' = E_{kin}^{1'} + E_{kin}^{2'} + V^{1'} + V^{2'}$$



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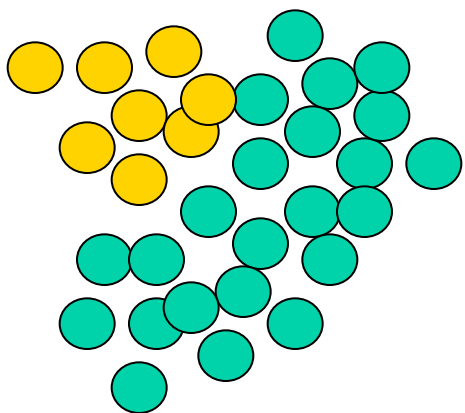
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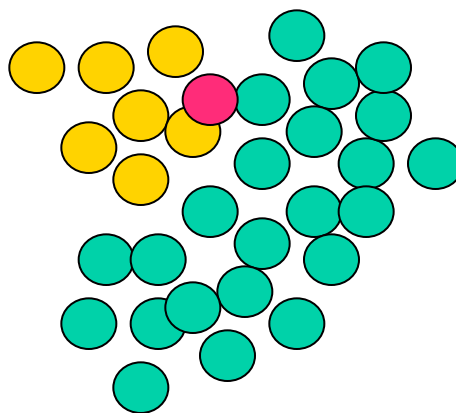
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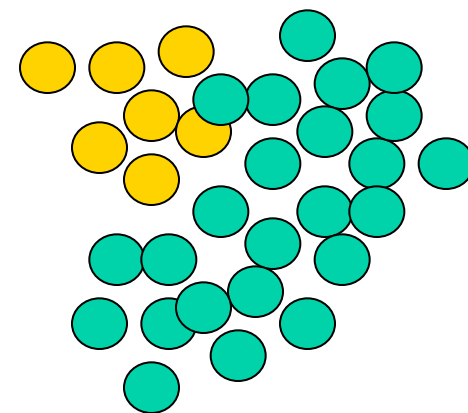
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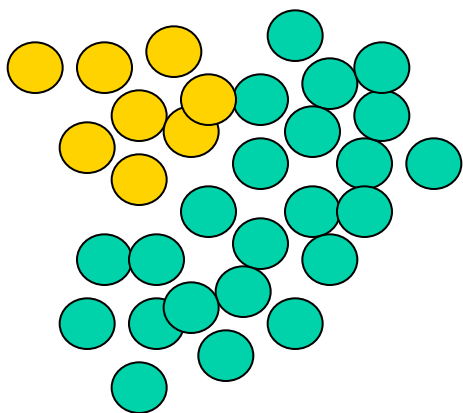
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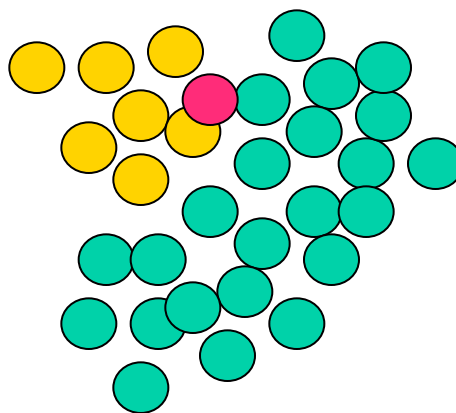
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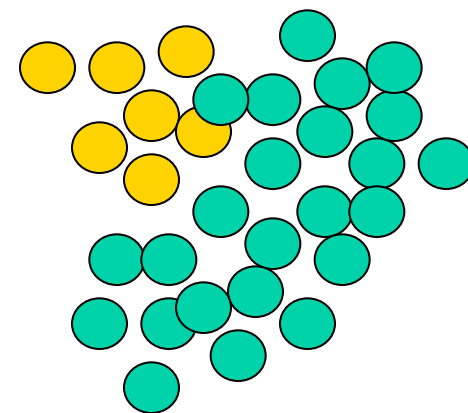
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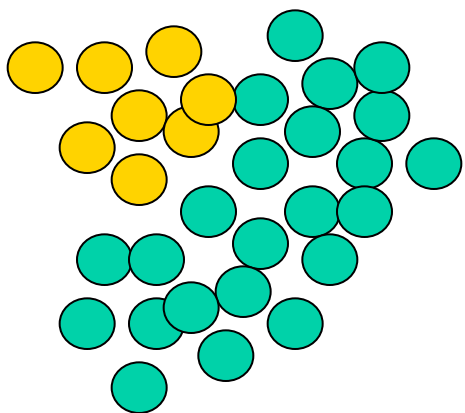
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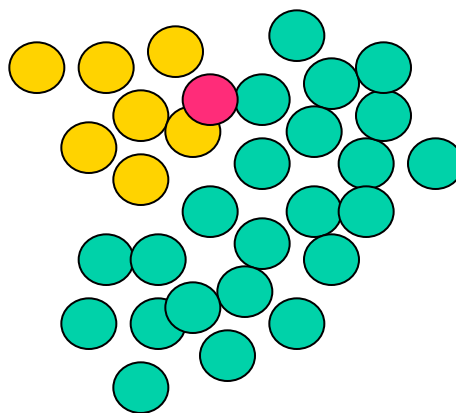
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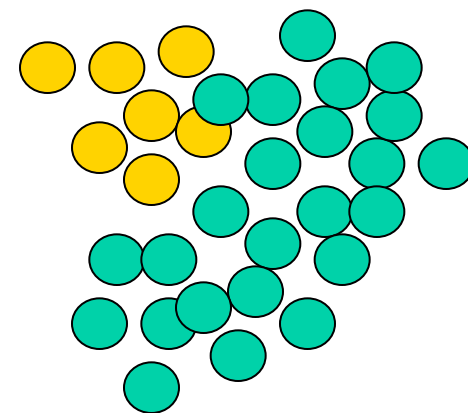
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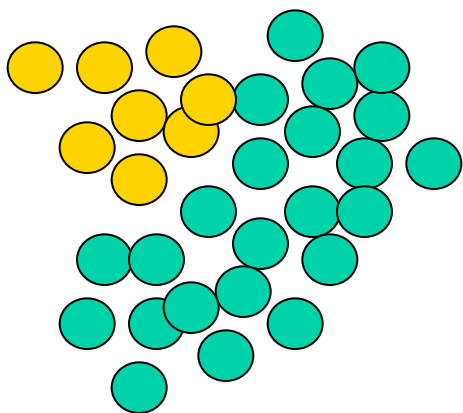
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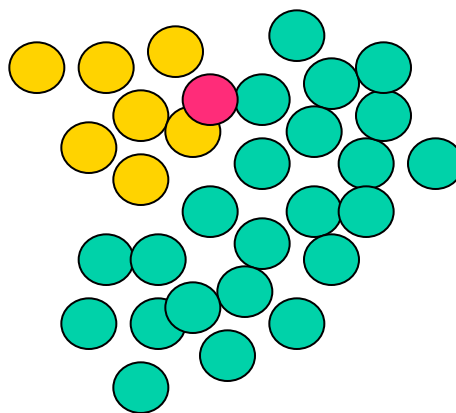
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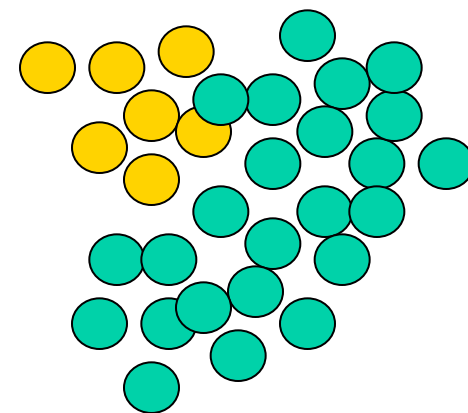
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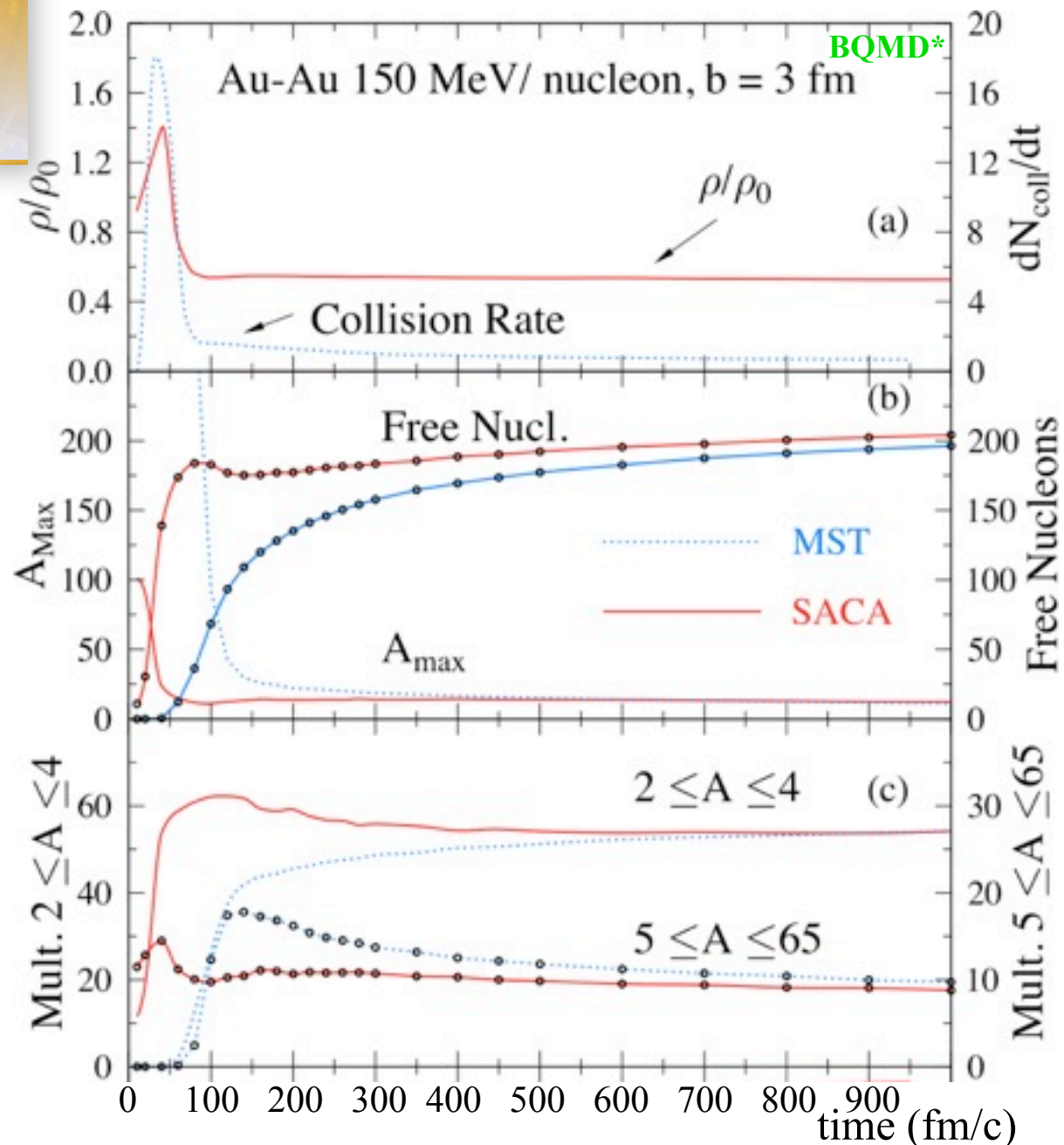
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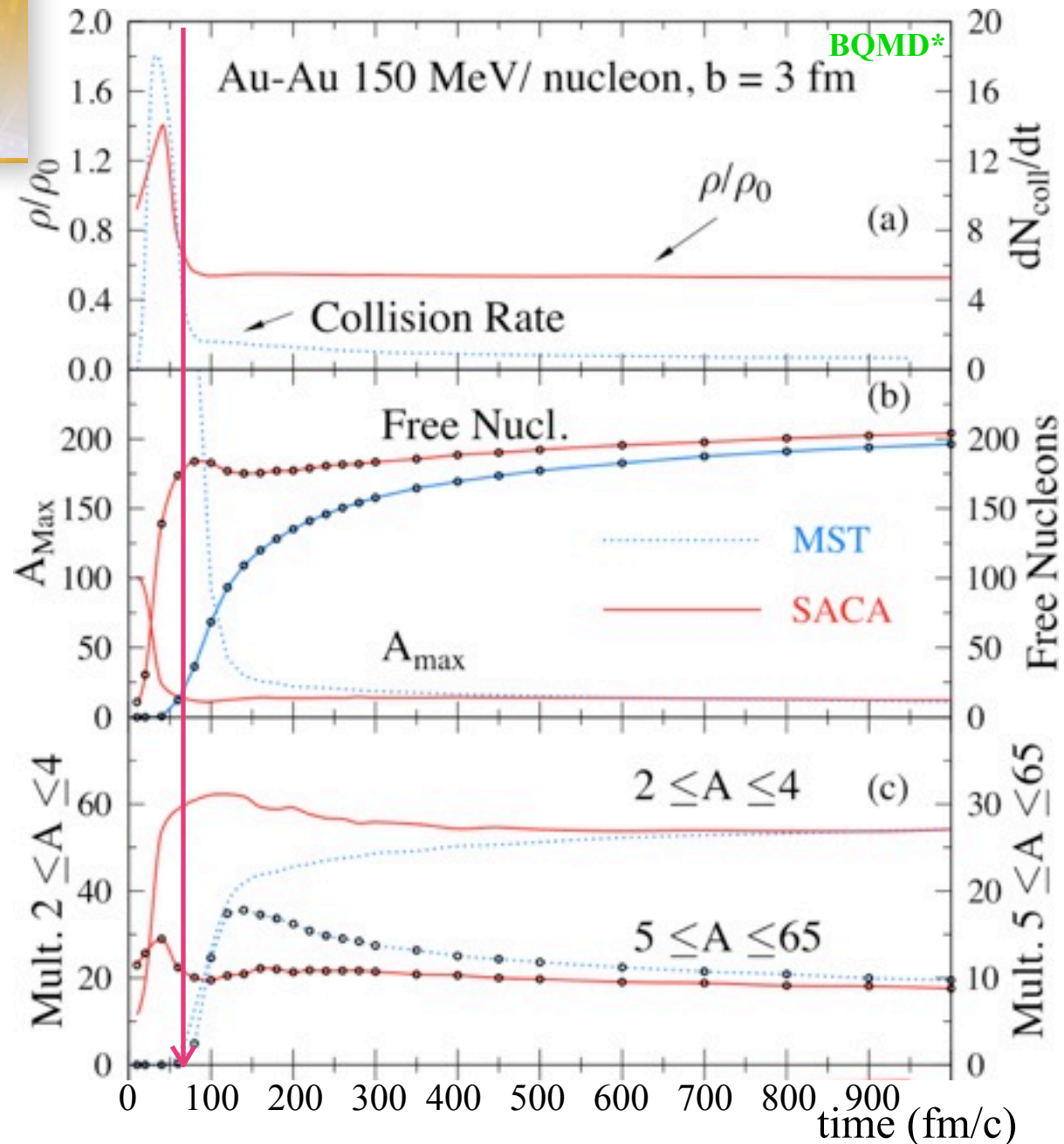
It leads automatically to the most bound configuration.

# SACA versus coalescence (Minimum Spanning Tree)



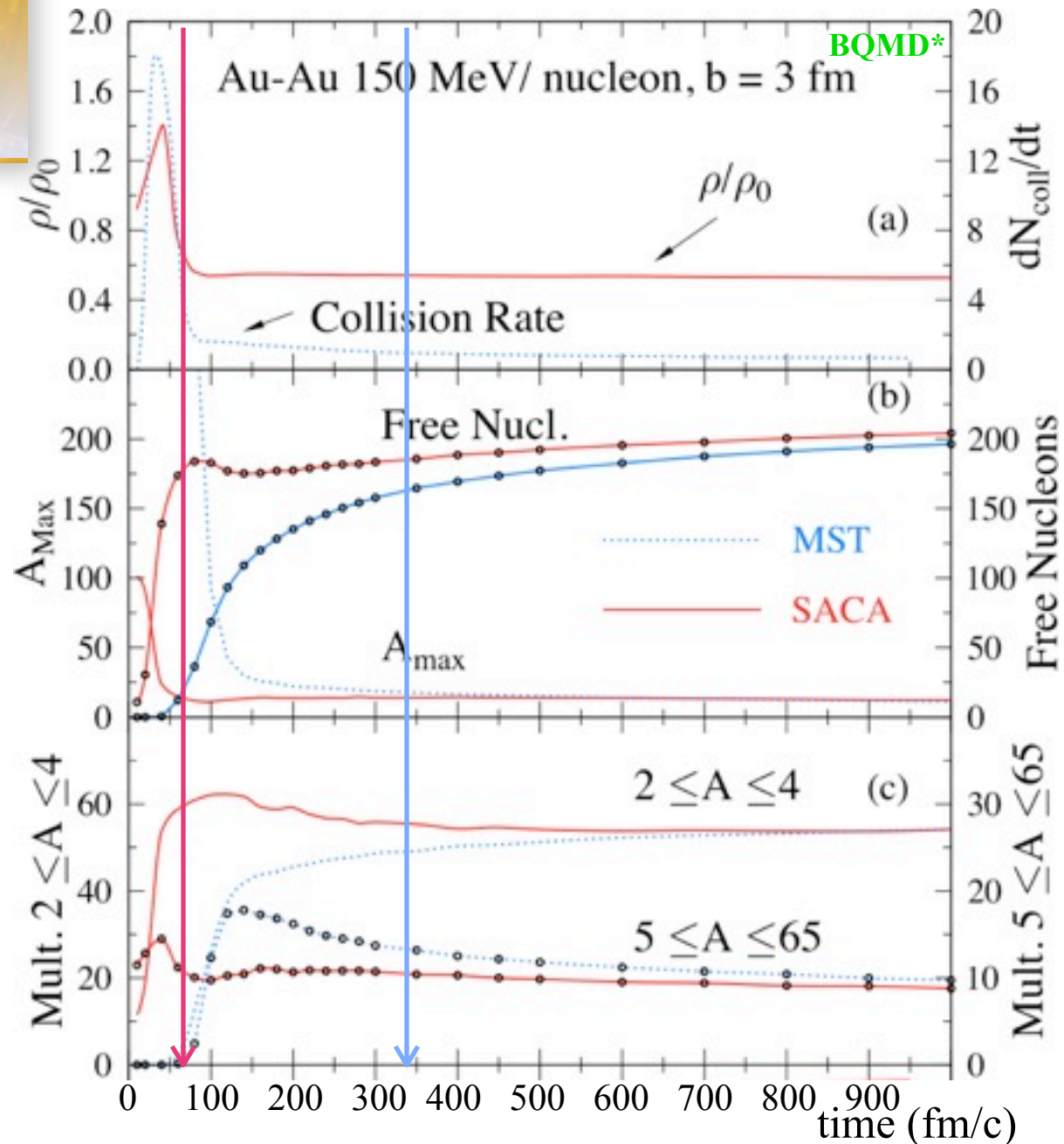
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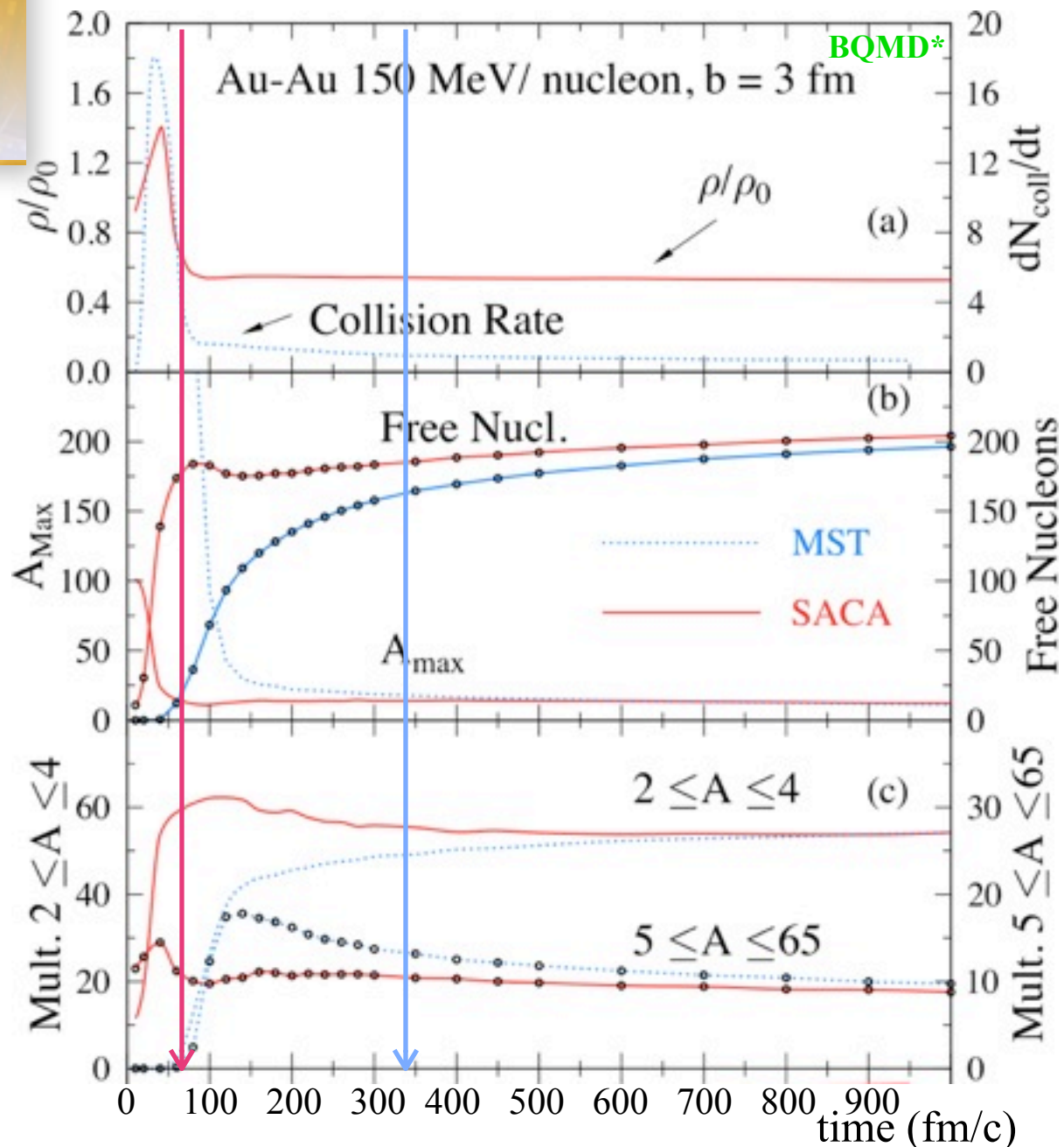
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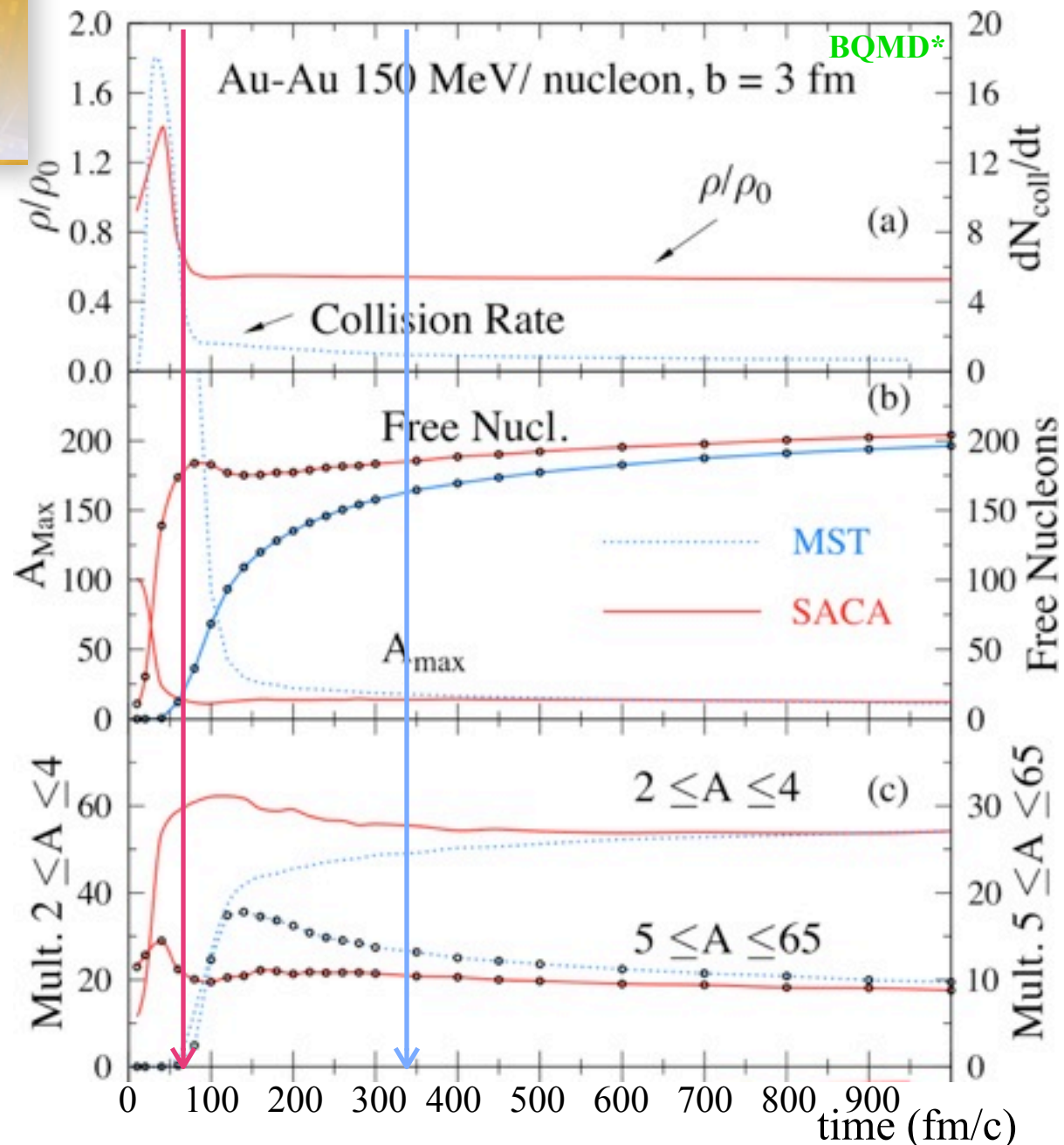


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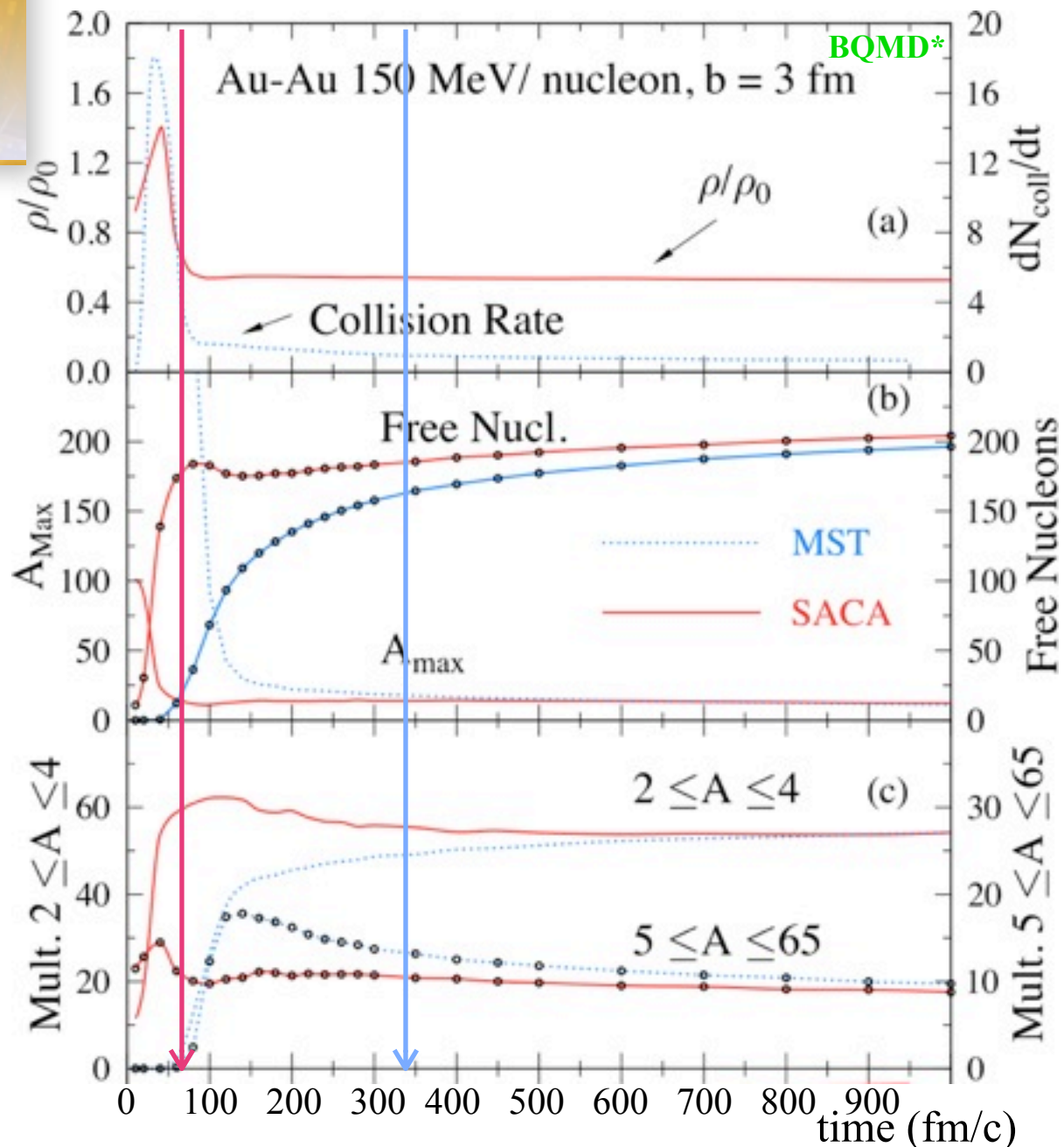
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▶ Advantage of SACA : the fragment partitions can reflect the early dynamical conditions (Coulomb, density, flow details, strangeness...).



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# Toward the isotope yields... IQMD + new SACA

SACA is applied here on the IQMD transport model\* calculations

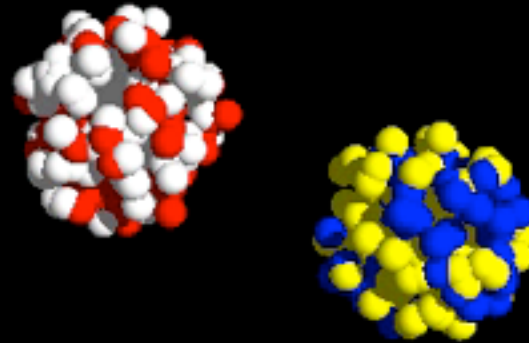
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Au+Au at 100 A.MeV -  $b=7$  fm



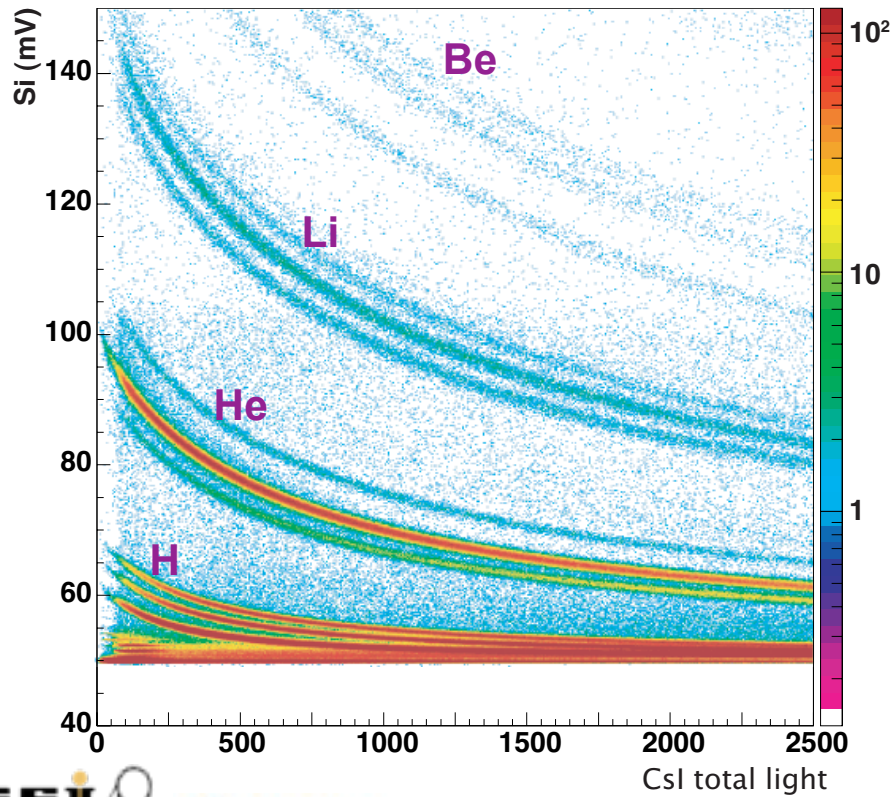
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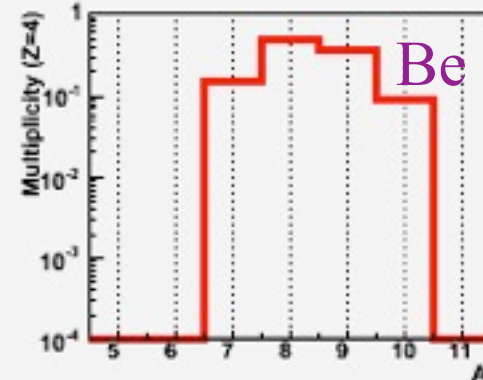
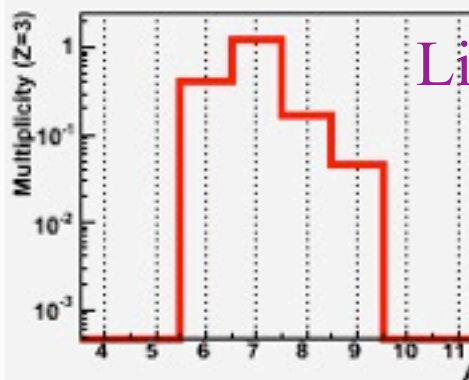
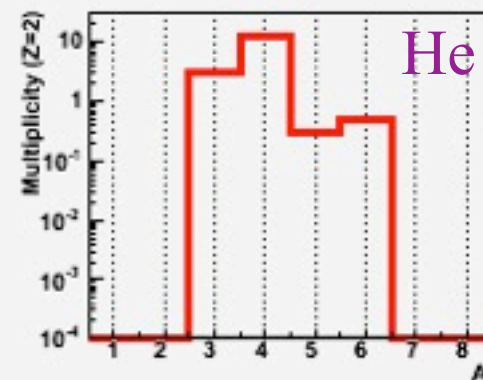
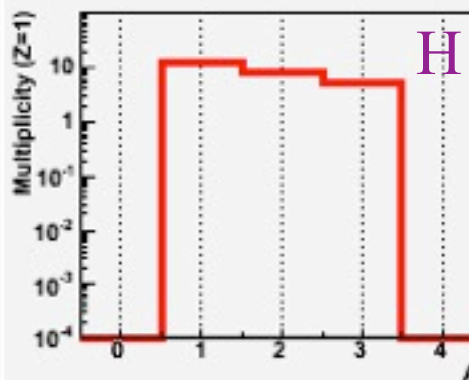
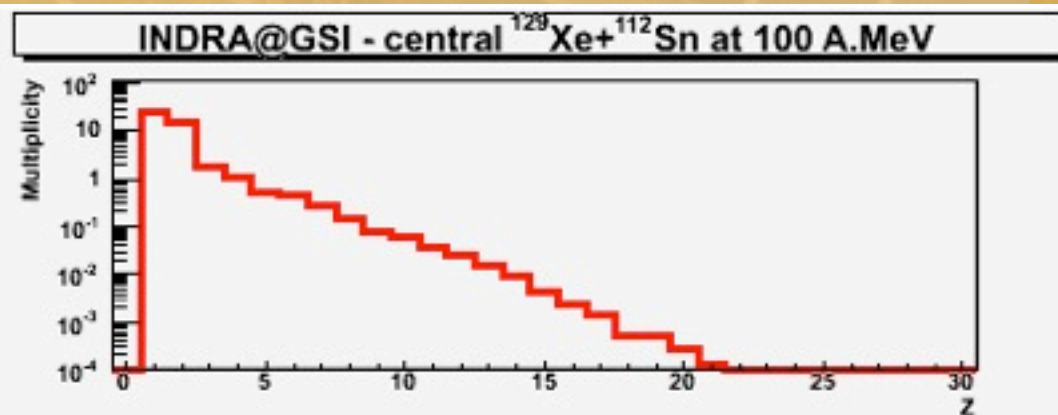
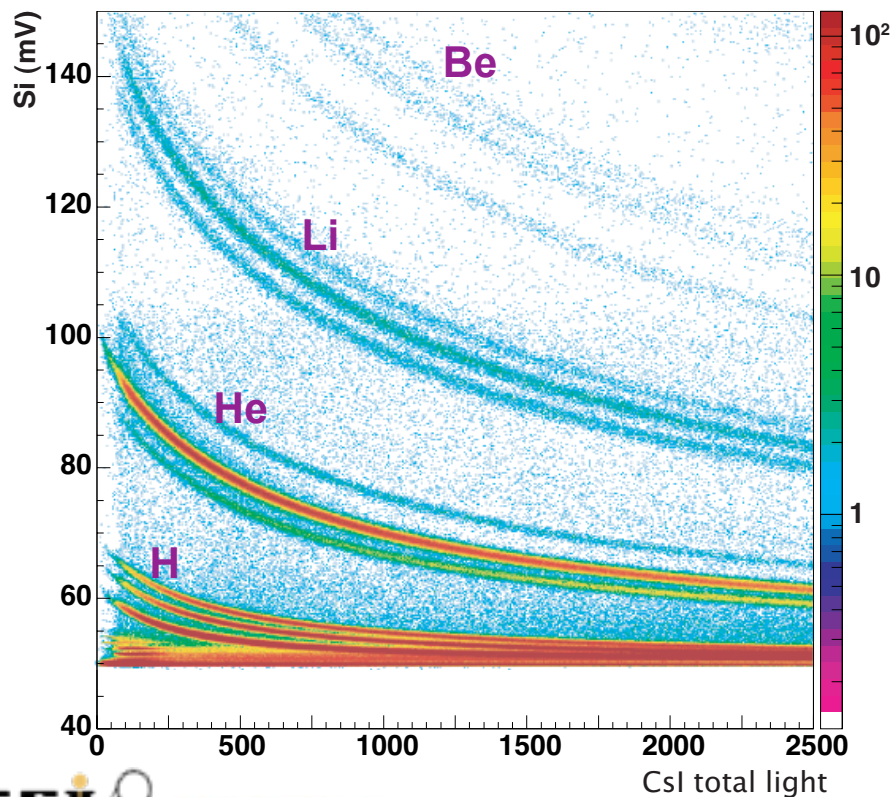
ALADiN - INDRA collaboration



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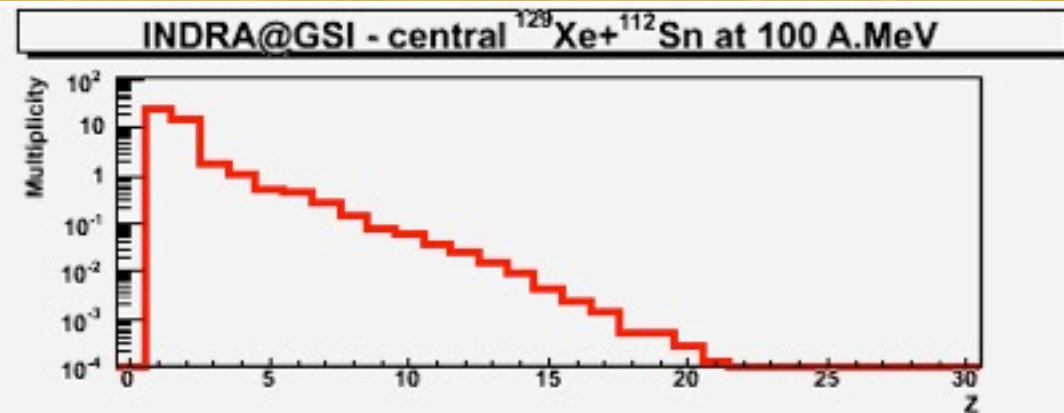
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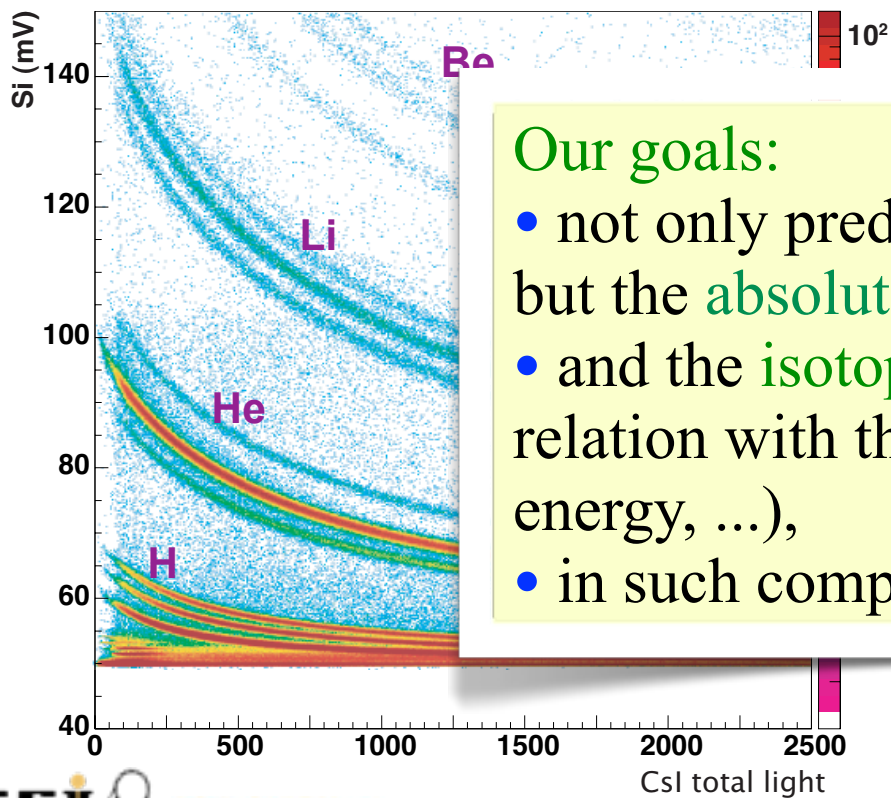


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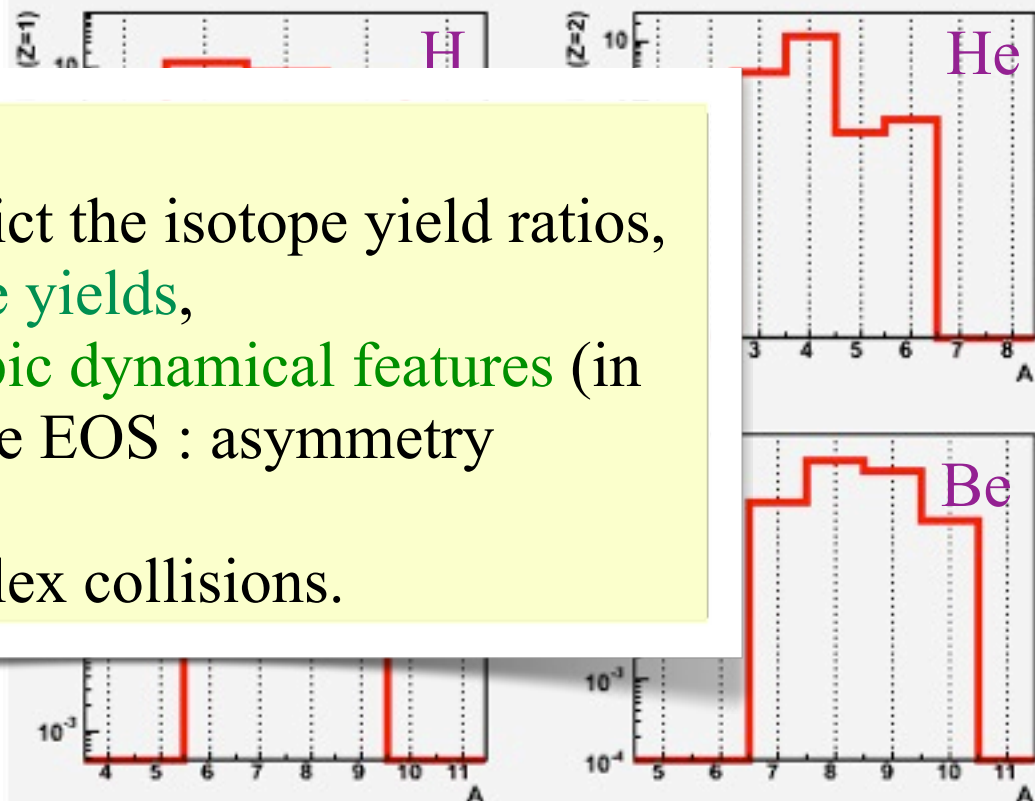


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Our goals:

- not only predict the isotope yield ratios, but the **absolute yields**,
- and the **isotopic dynamical features** (in relation with the EOS : asymmetry energy, ...),
- in such complex collisions.



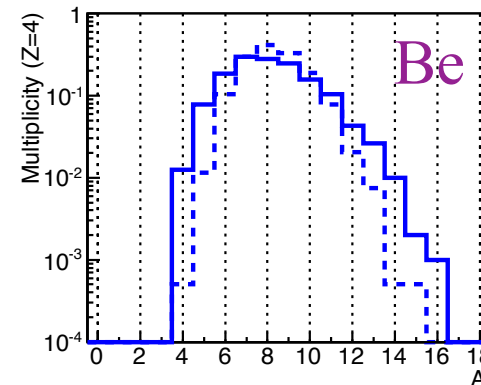
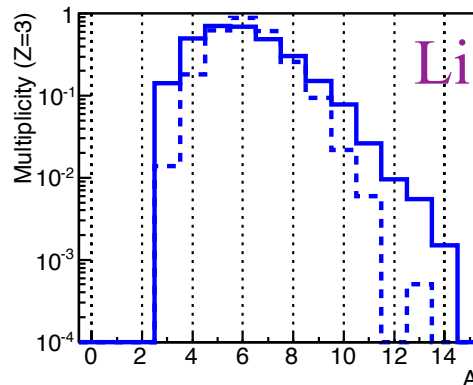
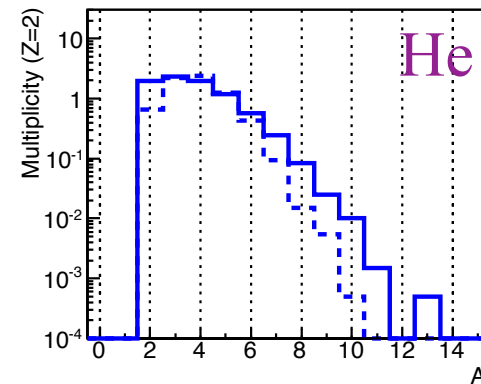
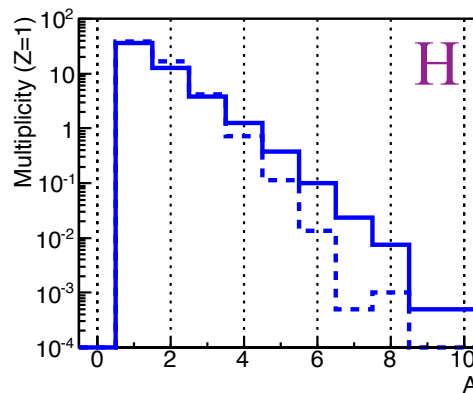
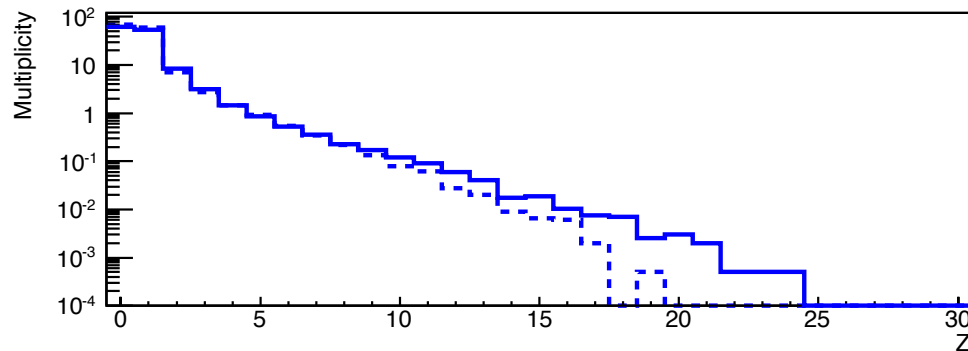
# SACA versus coalescence (Minimum Spanning Tree)

IQMD  $^{136}\text{Xe} + ^{112}\text{Sn}$  at 100 A.MeV,  $b=1$  fm,  $t_{\text{SACA}} = 60$  fm/c

SACA version:

----- MST only (200 fm/c)

—  $E_{\text{asy}}=0$ , no pairing

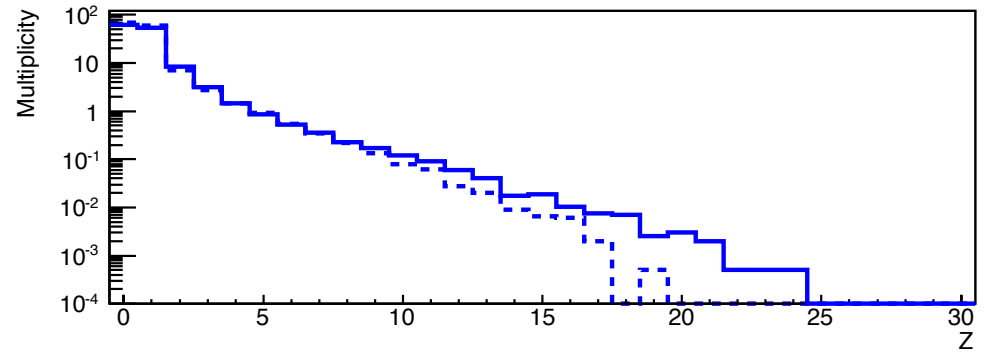


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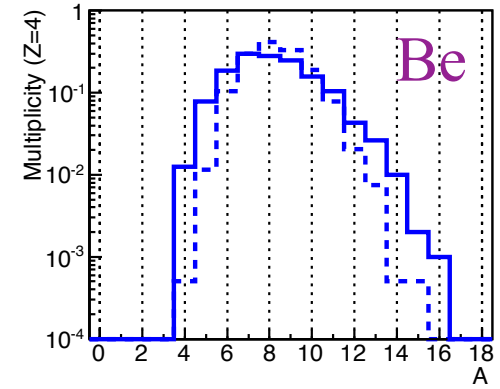
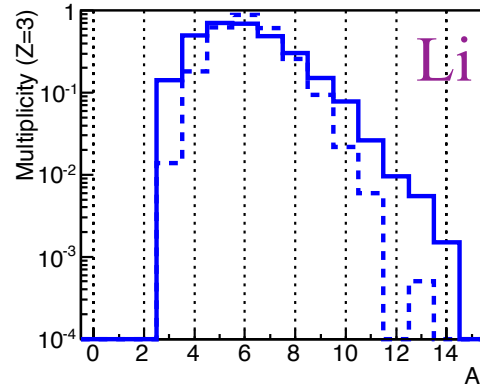
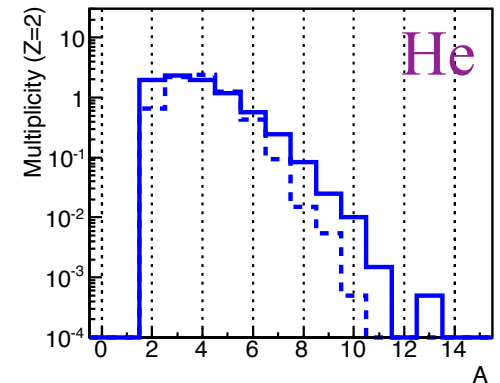
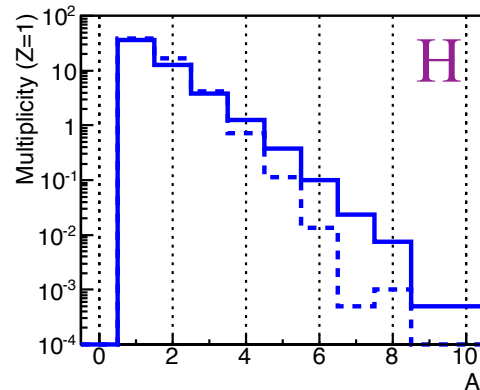
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At this stage, SACA contains as ingredients of the potential making the binding energy of the clusters :

- ① volume component: mean field (Skyrme, dominant)
- ② correction of surface effects: Yukawa

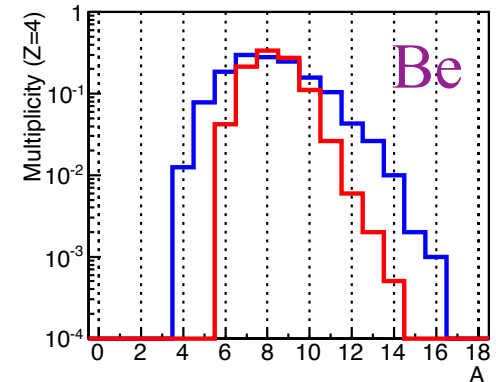
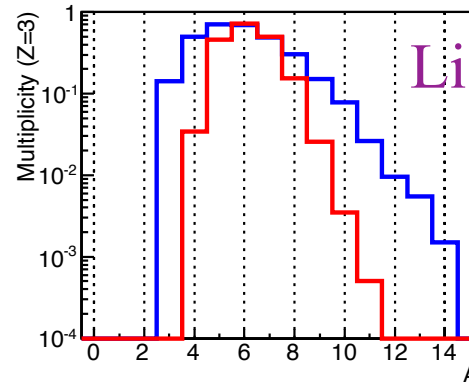
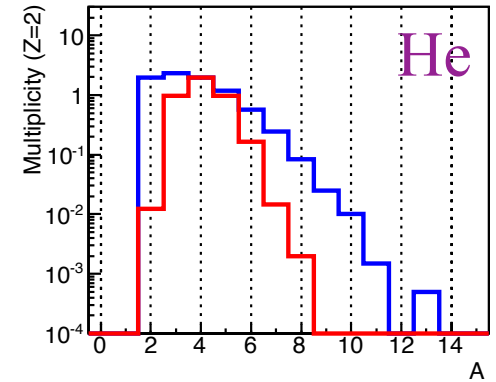
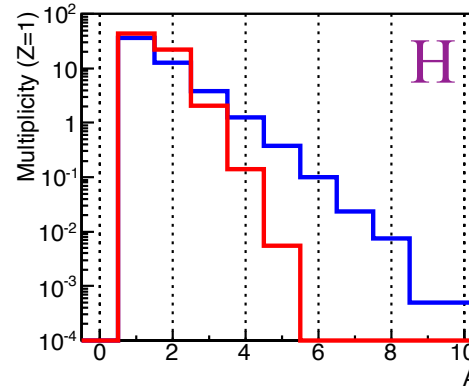
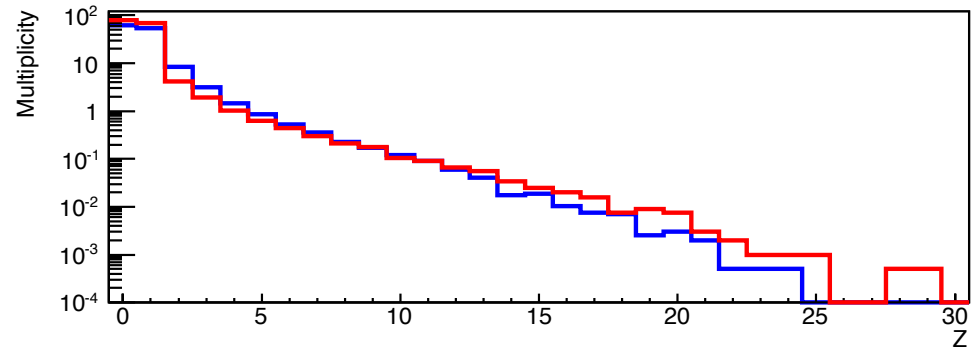


# SACA with asymmetry energy

IQMD  $^{136}\text{Xe} + ^{112}\text{Sn}$  at 100 A.MeV,  $b=1$  fm,  $t_{\text{SACA}}=60$  fm/c

SACA version:

- $E_{\text{asy}}=0$ , no pairing
- $E_{\text{asy}}=32$  MeV ( $\gamma=1$ ), no pairing

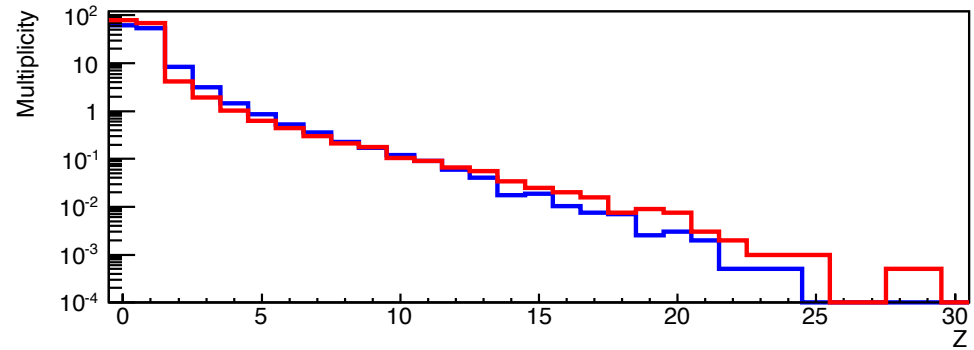


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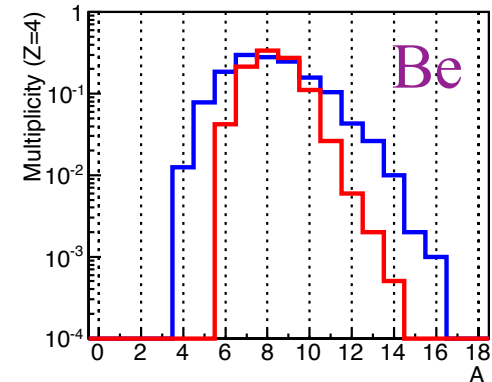
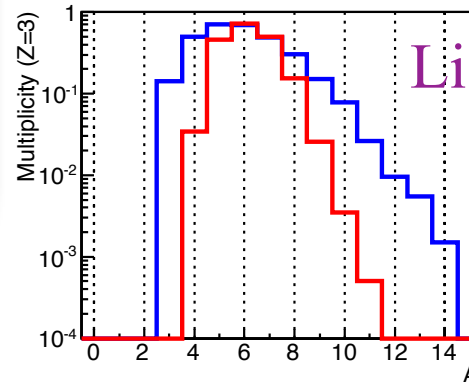
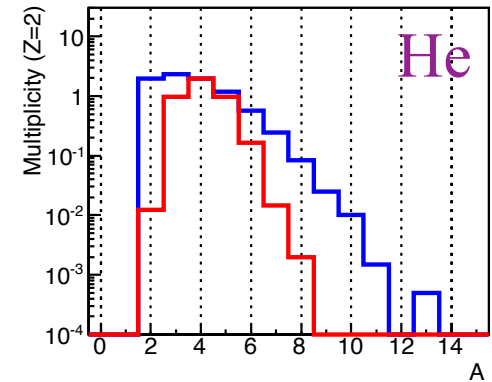
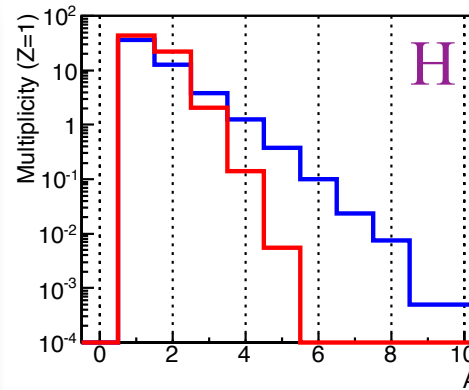


Here, in IQMD and SACA, we adopt the following asymmetry energy parametrisation:

$$E_{\text{asy}} = E_0 \cdot (\langle \rho_B \rangle / \rho_0)^{(\gamma-1)} \cdot (\langle \rho_n \rangle - \langle \rho_p \rangle) / \langle \rho_B \rangle$$

with  $E_0=32$  MeV,  $\gamma=1$  («stiff»)

- ➔ Z and A yields not strongly modified
- ➔ Isotope yields shrink onto the N=Z line
- ➔ Still not fully realistic: shell, odd-even effects (pairing) still absent.

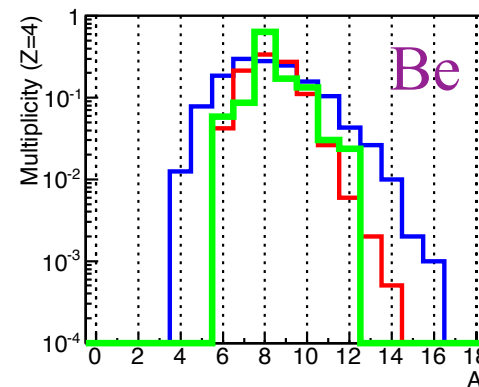
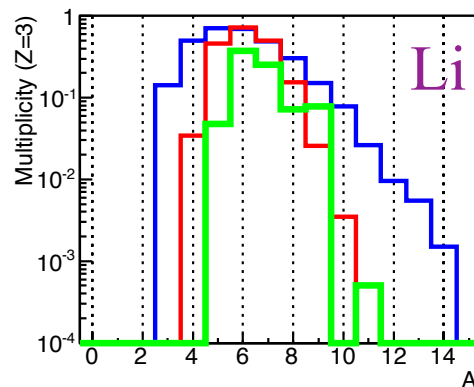
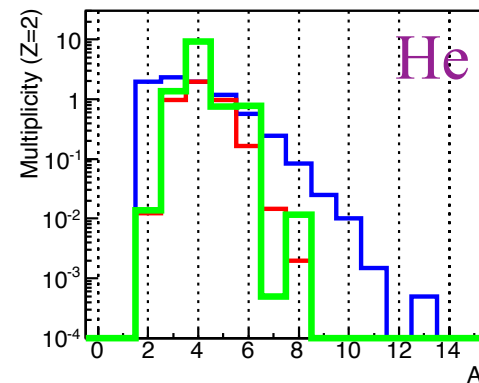
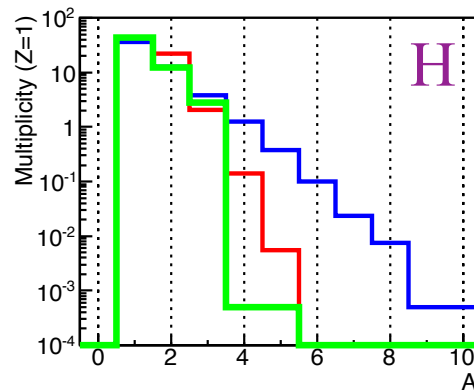
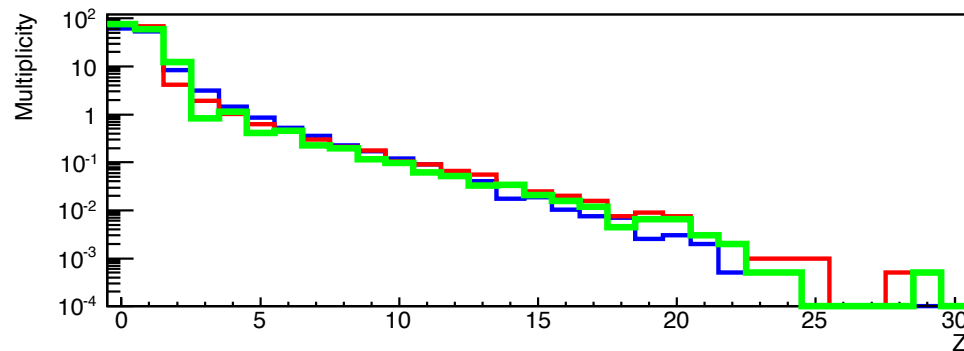


# SACA with asymmetry energy and pairing

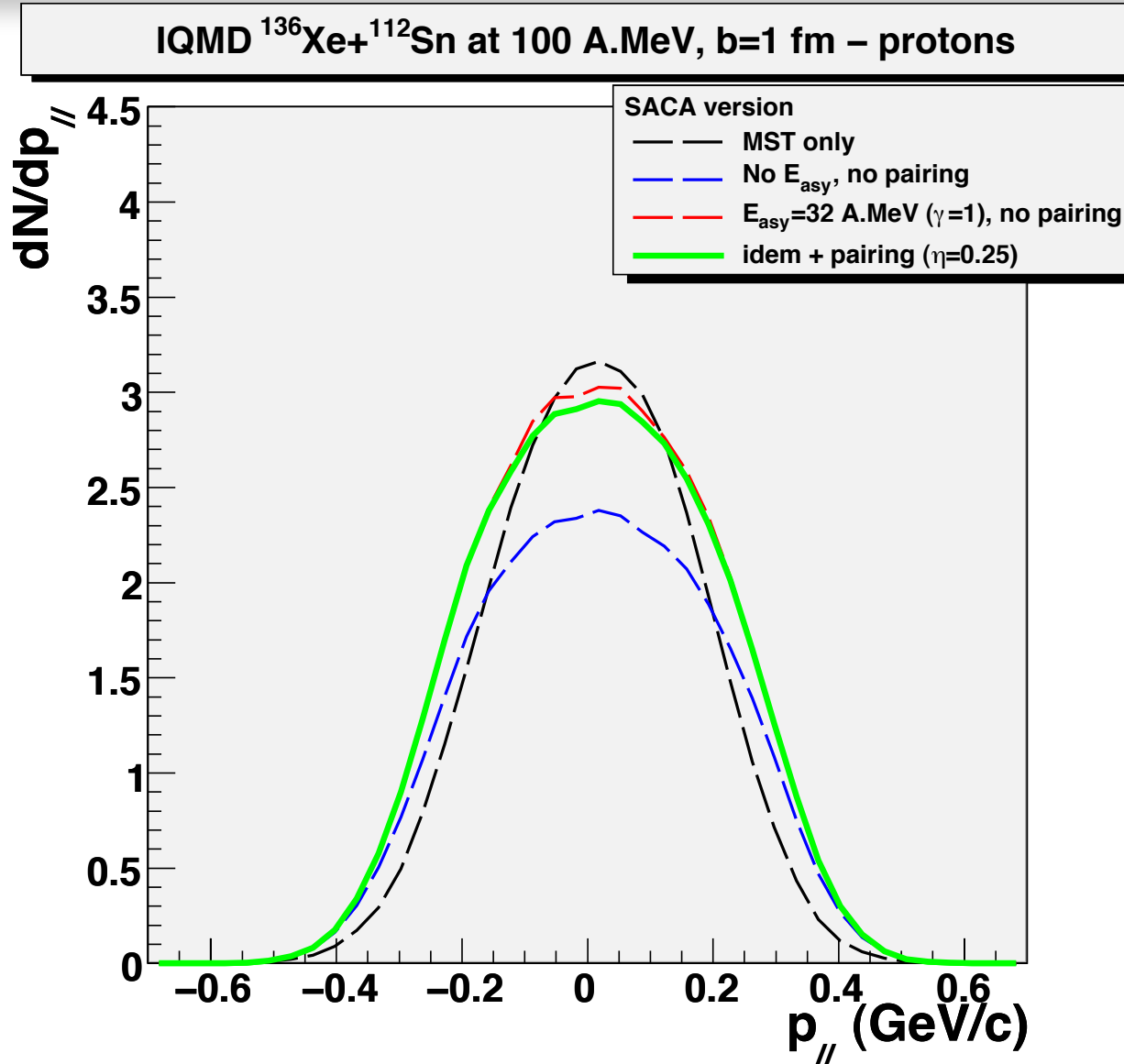
INDRA@GSI - central  $^{129}\text{Xe} + ^{112}\text{Sn}$  at 100 A.MeV

SACA version:

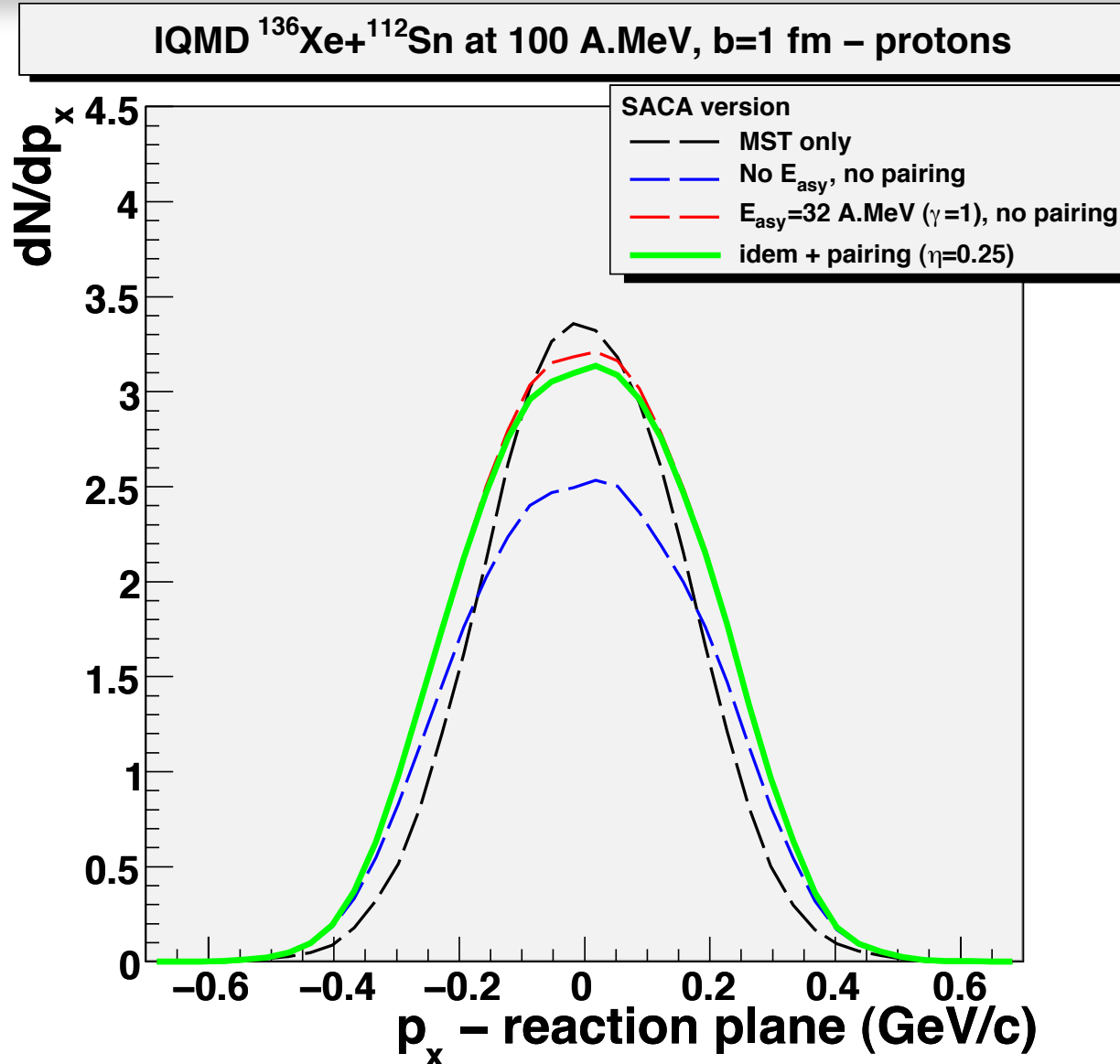
- $E_{\text{asy}}=0$ , no pairing
- $E_{\text{asy}}=32$  MeV ( $\gamma=1$ ), no pairing
- $E_{\text{asy}}=32$  MeV ( $\gamma=1$ ) +  $\eta_{\text{pairing}}=0.25$



# How the dynamical patterns of isotopes are affected

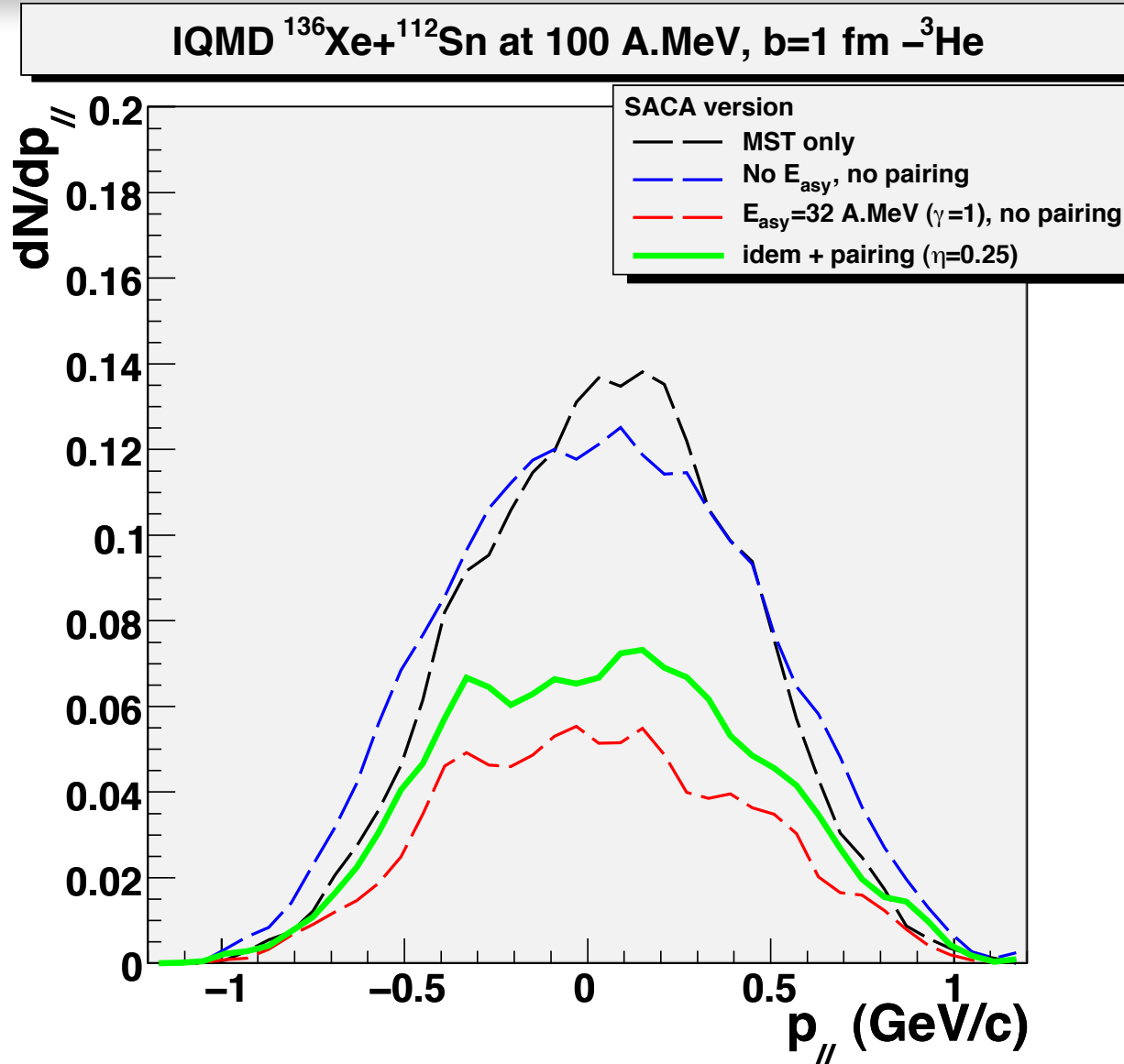


# How the dynamical patterns of isotopes are affected

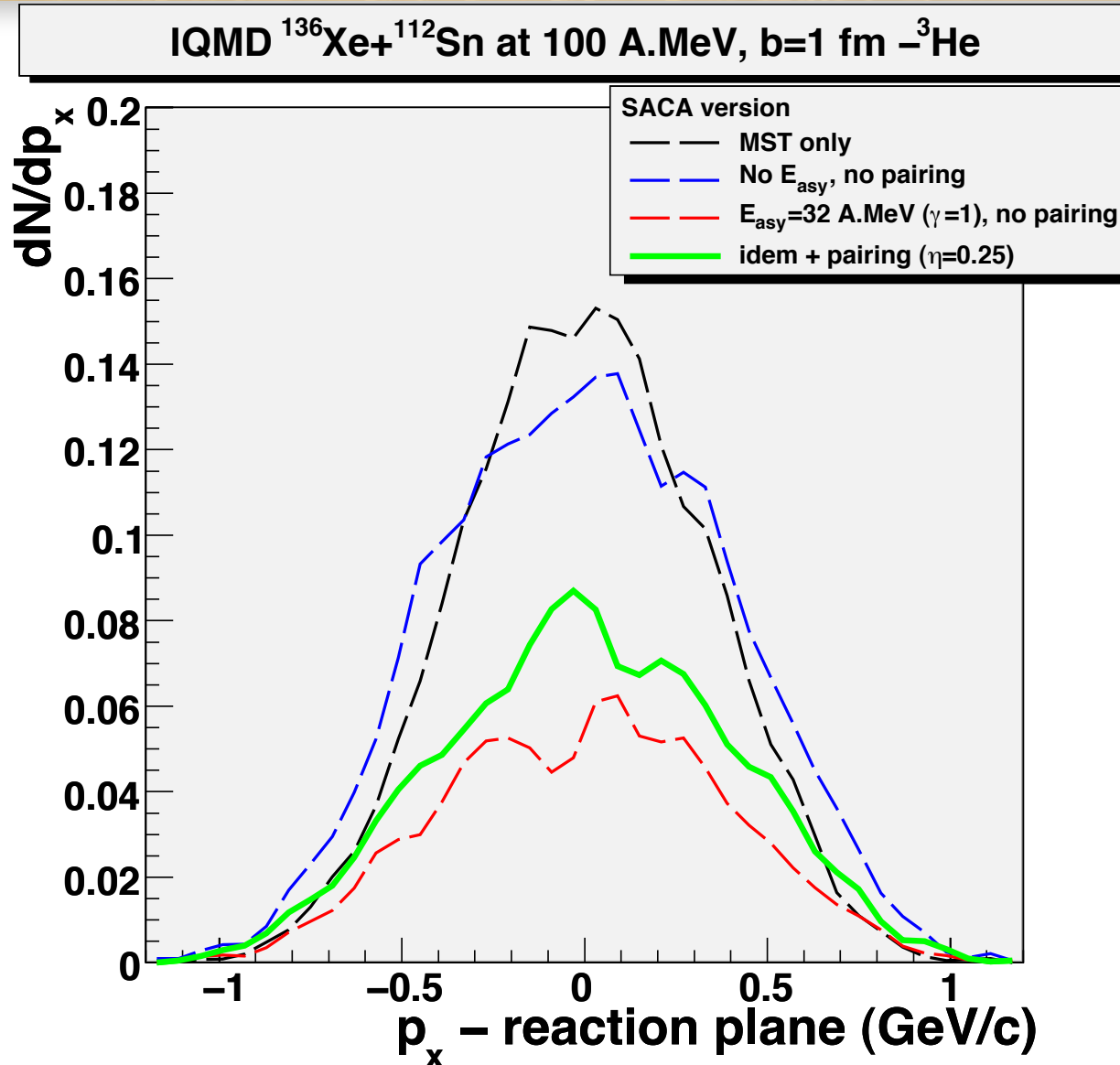




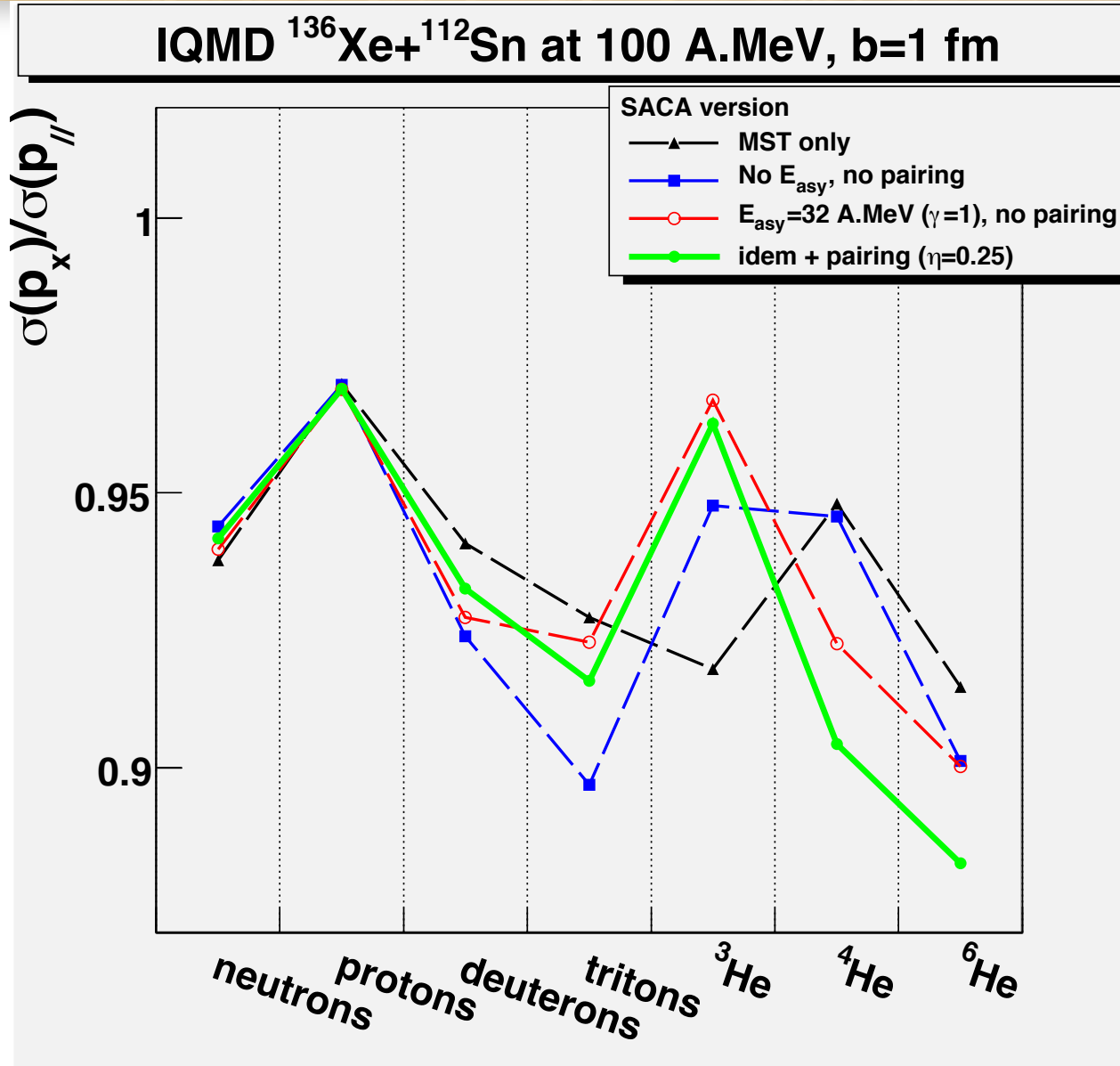
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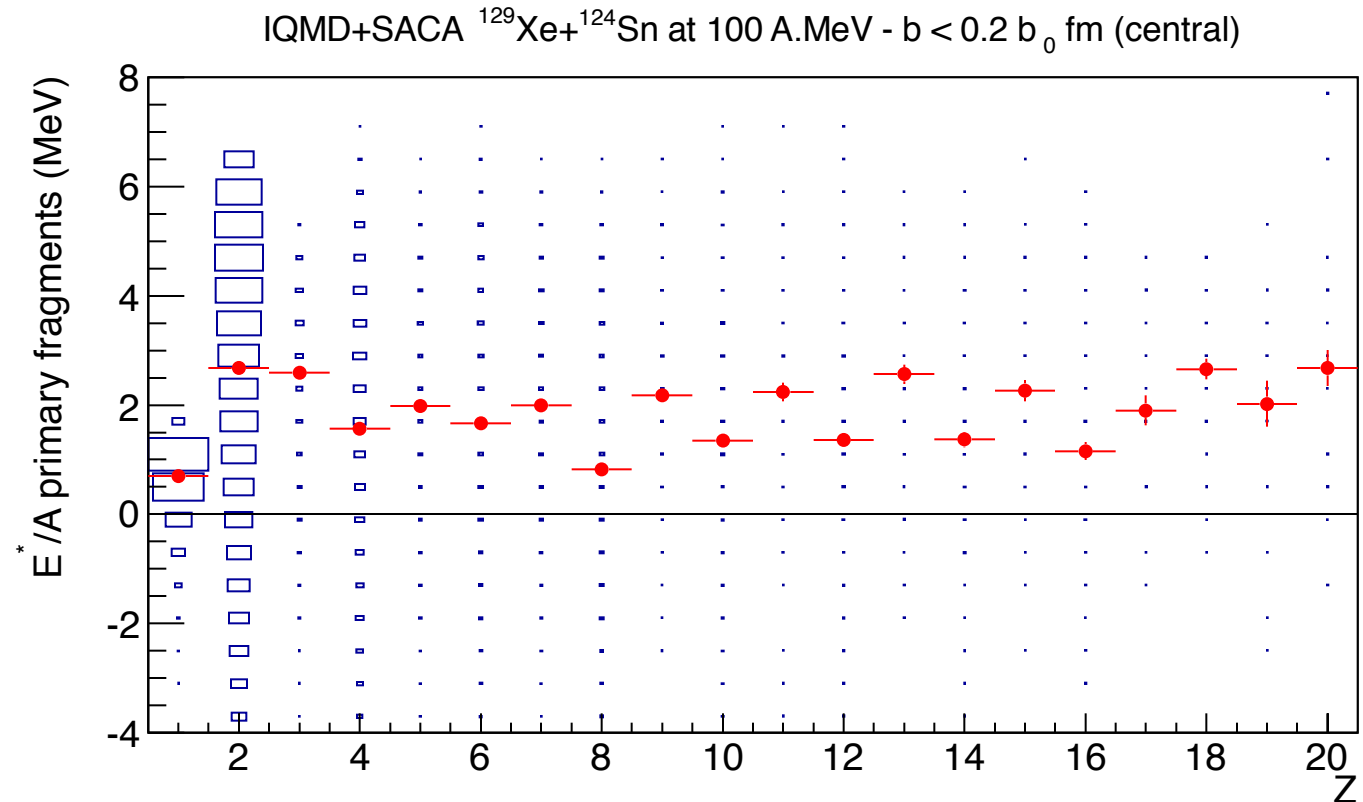
# Excitation energy of the primary fragments

$$E^* = E_{\text{g.s.}} - E_{\text{bind}}$$



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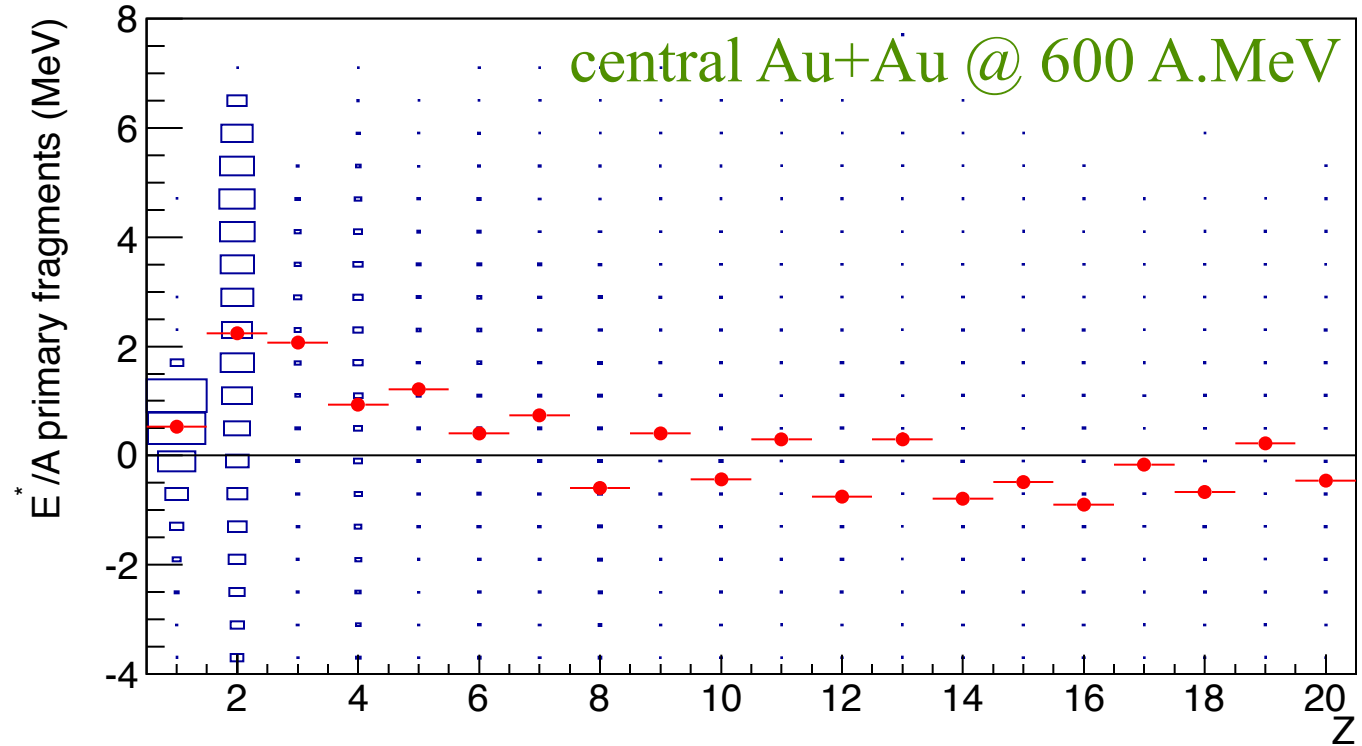
On Average, **2 A.MeV** of excitation energy. Corresponds to findings of S. Hudan et al. (INDRA collaboration), PRC 67, 064613 (2003).  
=> secondary decay (GEMINI) justified here.



# Excitation energy of the primary fragments

$$E^* = E_{\text{g.s.}} - E_{\text{bind}}$$

IQMD+SACA  $^{197}\text{Au} + ^{197}\text{Au}$  at 600 A.MeV -  $b < 0.2 b_0$  fm (central)



At relativistic energies, in the participant-spectator regime, heavy primary clusters are produced colder on average.



# What can we learn from the isotope yields regarding the asymmetry energy?

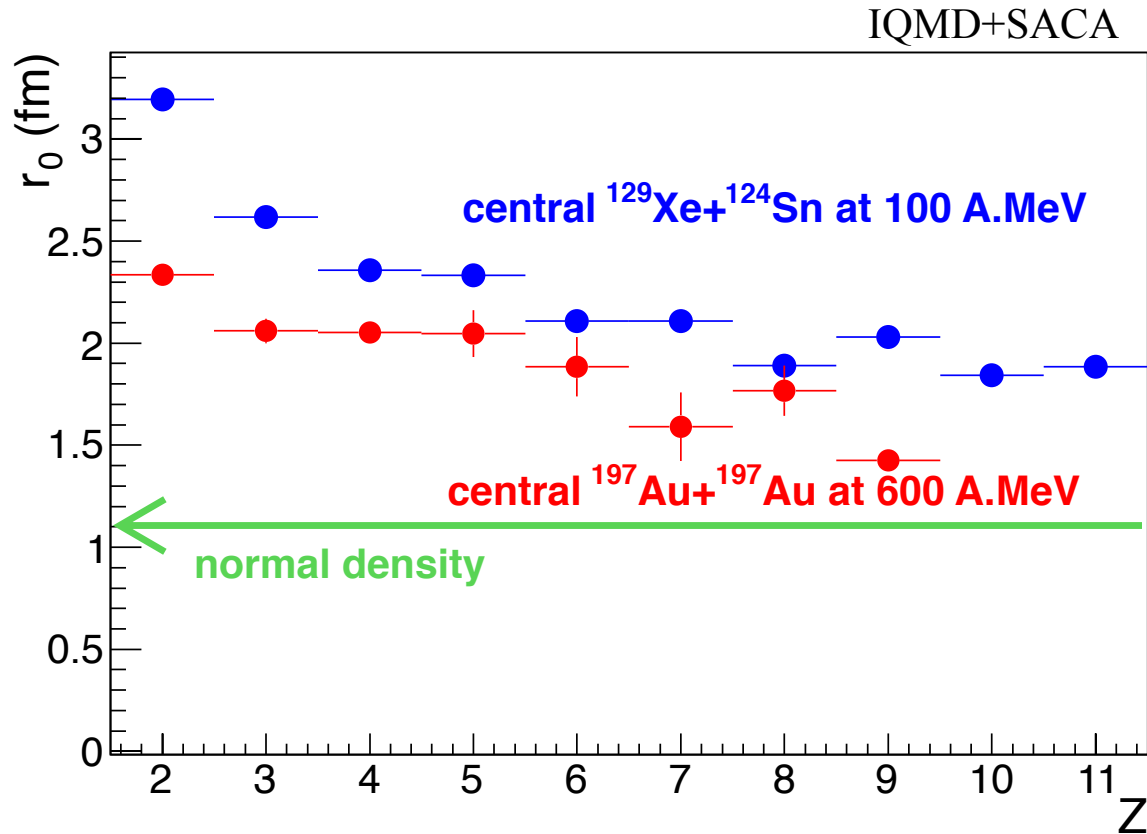
Mean radius of primary clusters

IQMD+SACA



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## Mean radius of primary clusters

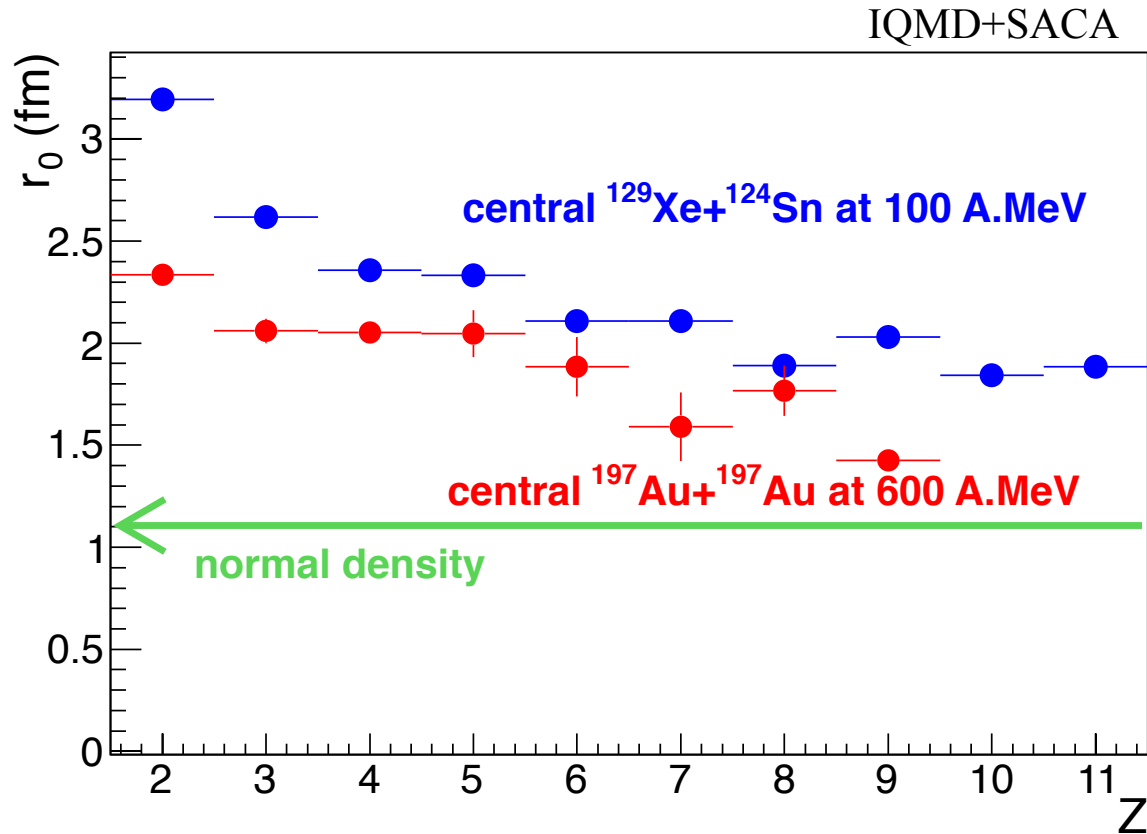






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## Mean radius of primary clusters

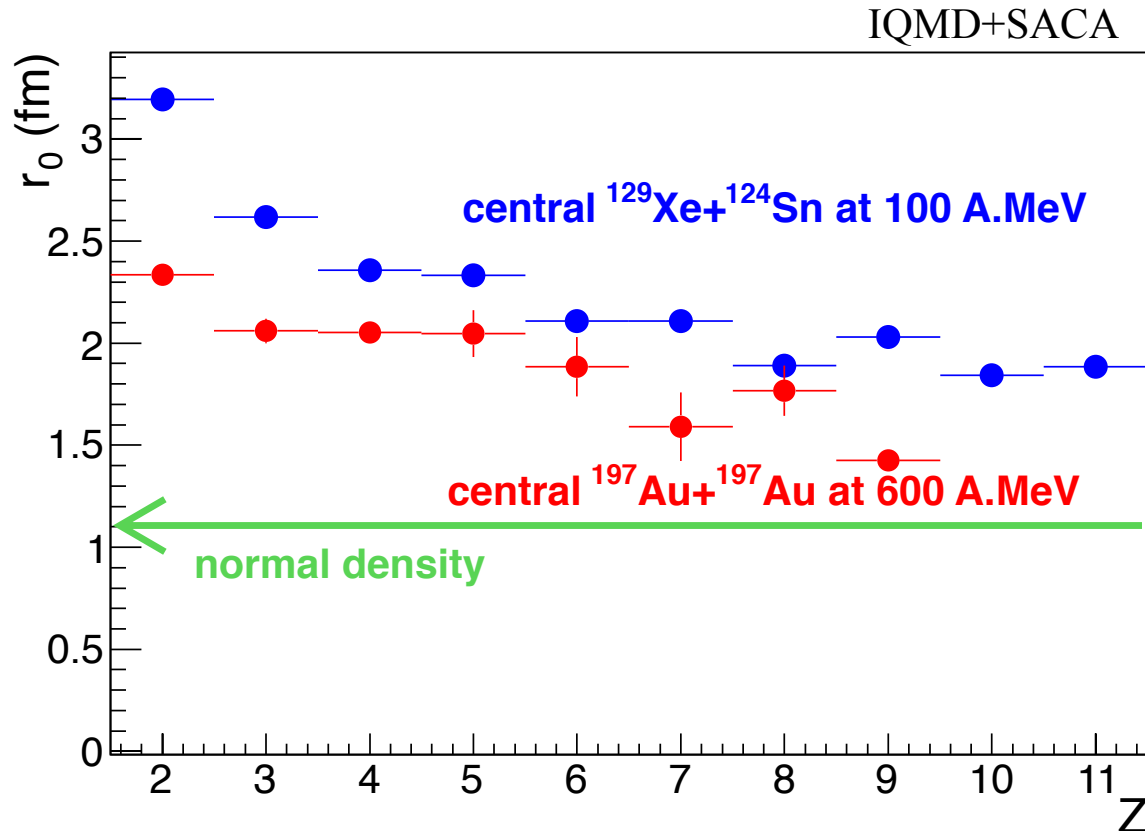


Though the medium is dense at this early stage, the dense clusters are disfavoured, because they would correspond to nucleons flowing against each other, hence with too high relative momenta to make a cluster.



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## Mean radius of primary clusters



Though the medium is dense at this early stage, the dense clusters are disfavoured, because they would correspond to nucleons flowing against each other, hence with too high relative momenta to make a cluster.

=> The isotope yields can only inform on the low density dependence on the asymmetry energy.



# Asymmetry energy influence versus system energy



# Asymmetry energy influence versus system energy

*W. Reisdorf and the FOPI Collaboration*

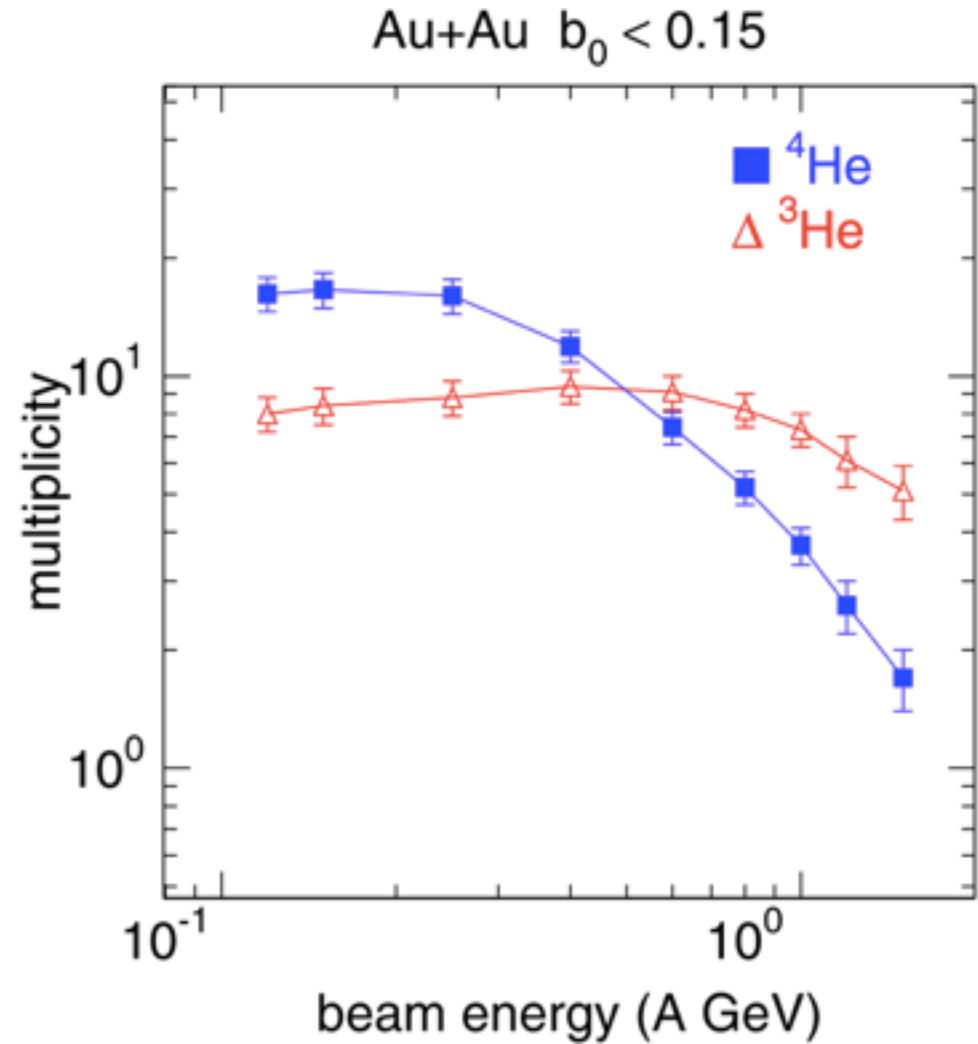
*Nuclear Physics A 848 (2010) 366–427*



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W. Reisdorf and the FOPI Collaboration

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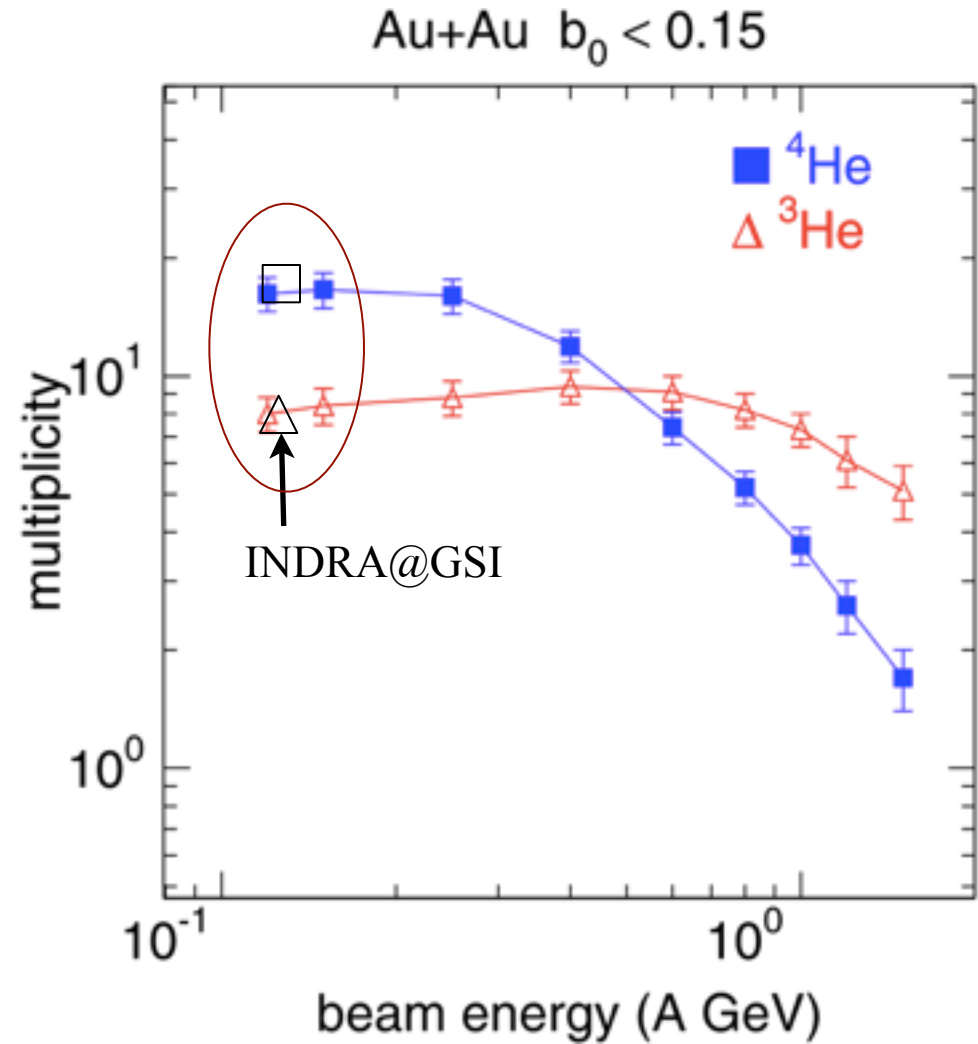




# Asymmetry energy influence versus system energy

W. Reisdorf and the FOPI Collaboration

Nuclear Physics A 848 (2010) 366–427

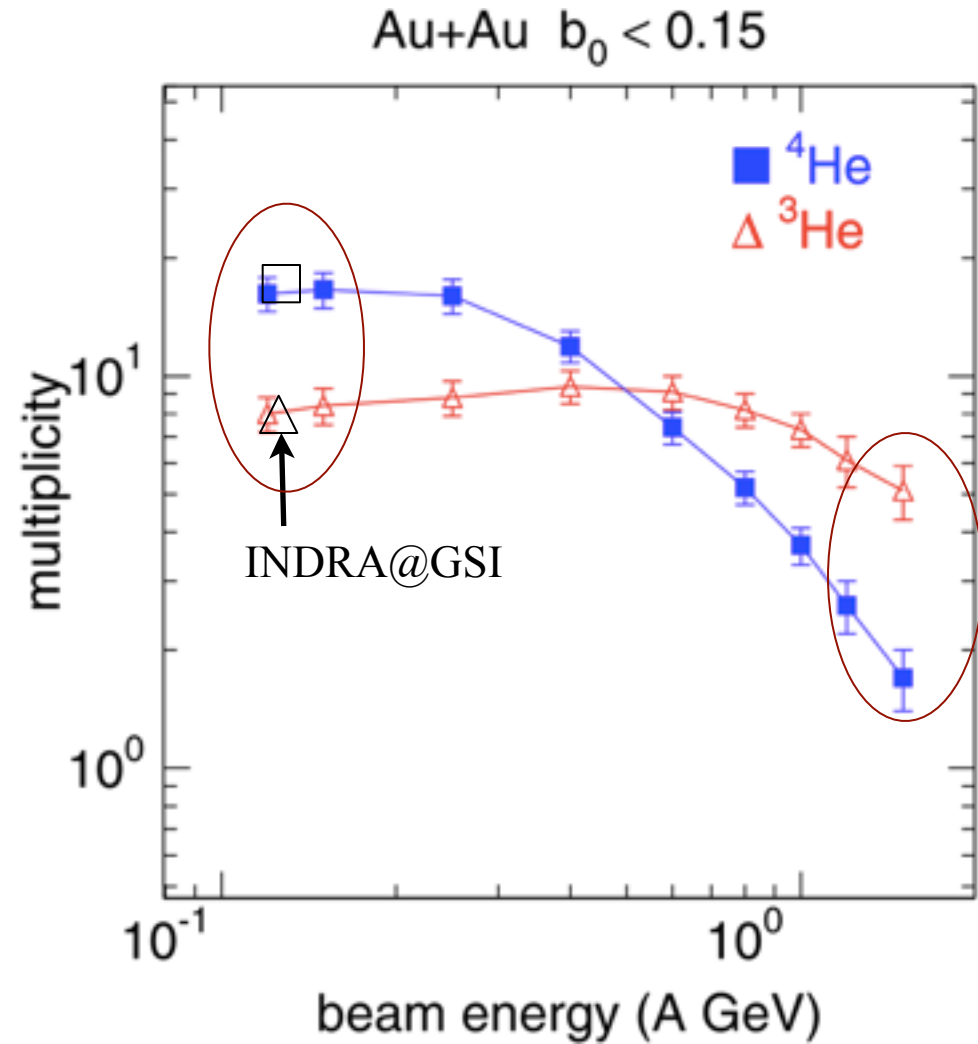




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Nuclear Physics A 848 (2010) 366–427



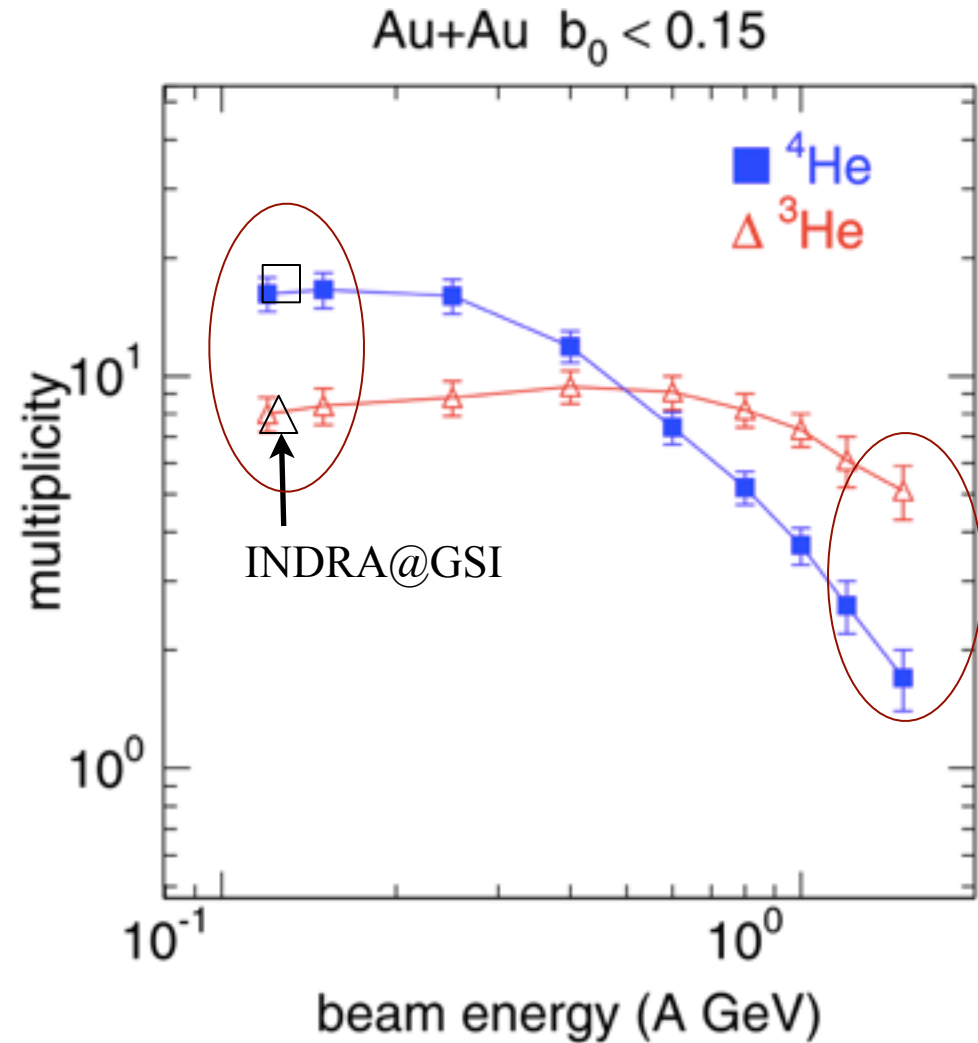


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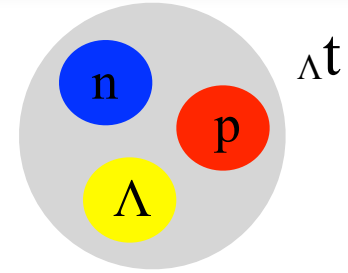
=> At high energy, the asymmetry energy effect on clusters seems to vanish. Timescale effect? Non-linear dependence on the density?







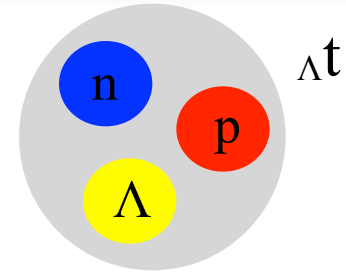
# Another application of SACRA : hypernuclei production





## Another application of SACRA : hypernuclei production

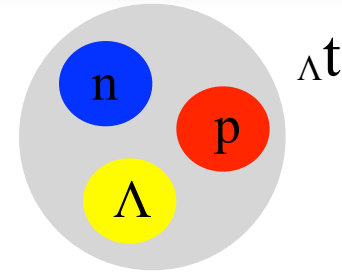
A hypernucleus is a nucleus which contains at least one hyperon ( $\Lambda$ (uds), ...) in addition to nucleons.





## Another application of SACA : hypernuclei production

A hypernucleus is a nucleus which contains at least one hyperon ( $\Lambda$ (uds), ...) in addition to nucleons.



Extending SACA for clusterising hadrons with hyperons (lambdas,...) for making **hypernuclei** is straightforward:

- ❖ one replaces  $V_{n-p}$  by  $V_{\Lambda-p}$  and  $V_{n-n}$  by  $V_{\Lambda-n}$
- ❖ and applies with these modifications the SACA algorithm.

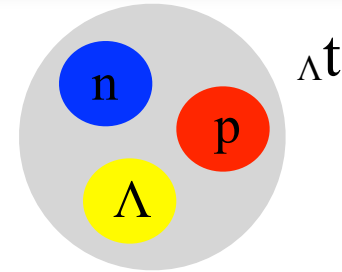
If in a final fragment there is a lambda, a hypernucleus should be created.

As a first approach, we have adopted  $V_{\Lambda-N} = 2/3 V_{n-N}$  ; further refinements are possible.



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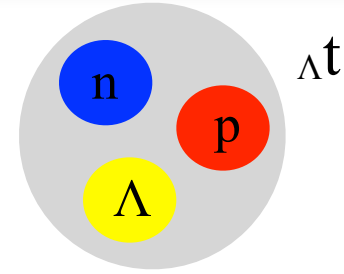
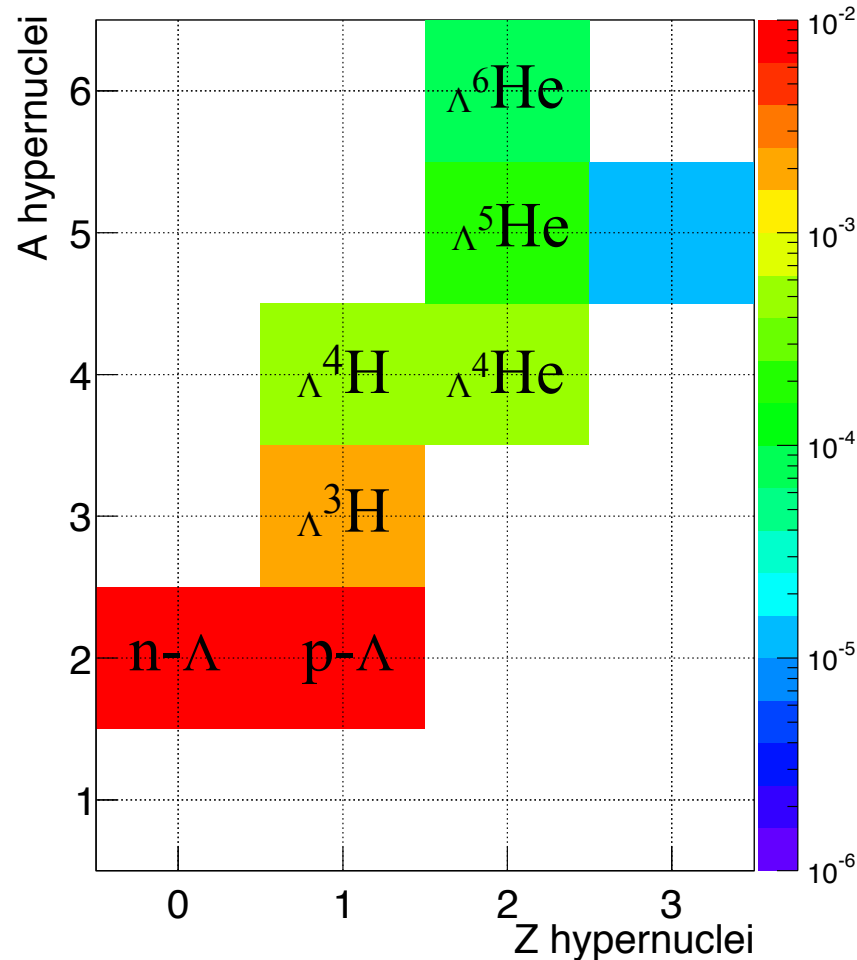
Many ways of producing lambdas:  $K N \rightarrow \Lambda \pi$ ,  $\pi^+ n \rightarrow \Lambda K^+$ ,  $\pi^- p \rightarrow \Lambda K_0$ ,  $p p \rightarrow \Lambda X$   
 $\Rightarrow$  influence of the EOS, in medium-properties, etc.



# Another application of SACA : hypernuclei production

IQMD+SACA  
 $^{58}\text{Ni}+^{58}\text{Ni}$   
 at  
 1.91 A.GeV  
 ( $b < 6 \text{ fm}$ ) -  
 $t_{\text{cluster.}} = 20 \text{ fm/c}$

IQMD+SACA  $^{58}\text{Ni}+^{58}\text{Ni}$  at 1.93 A.GeV ( $b < 6 \text{ fm}$ ,  $t_{\text{cluster.}} = 20 \text{ fm/c}$ ) - soft no mdi, kaon pot.



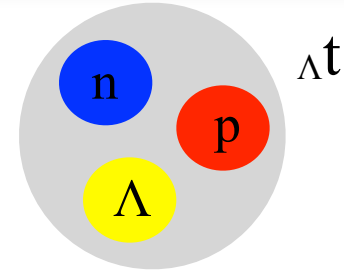
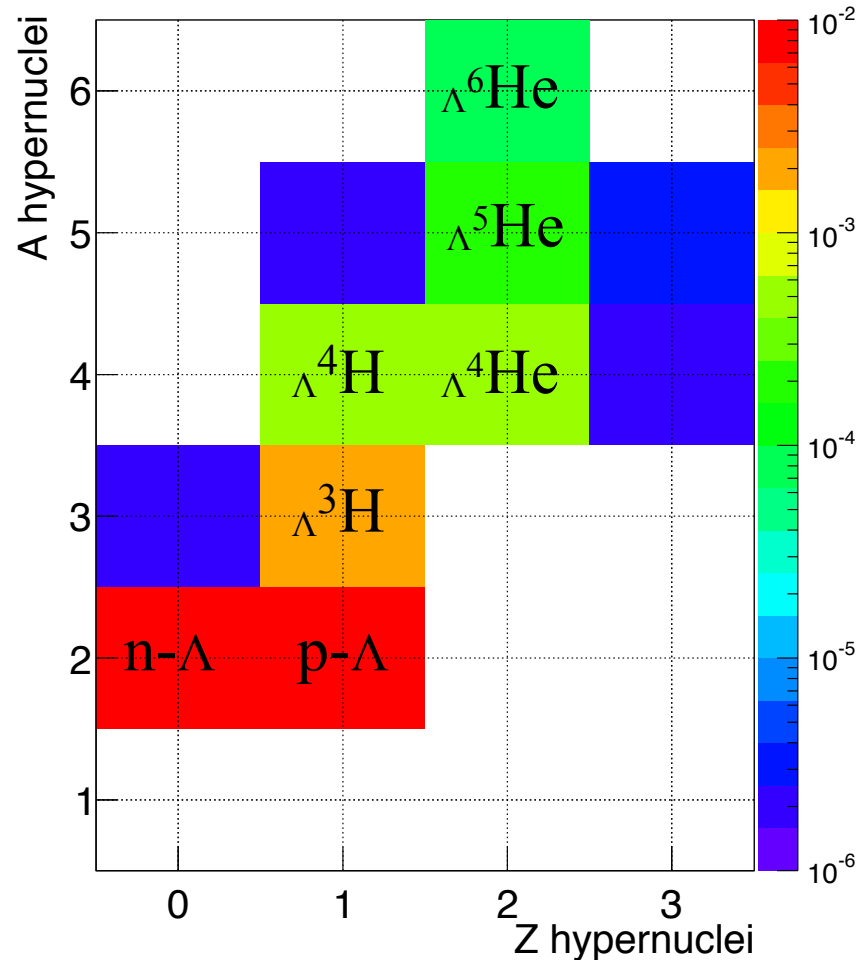
Soft EOS  
 no m.d.i.  
 with Kaon pot.



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Soft EOS  
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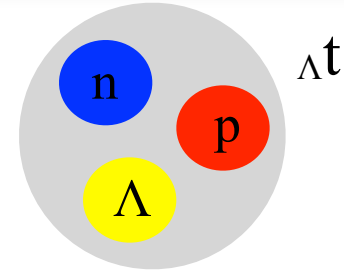
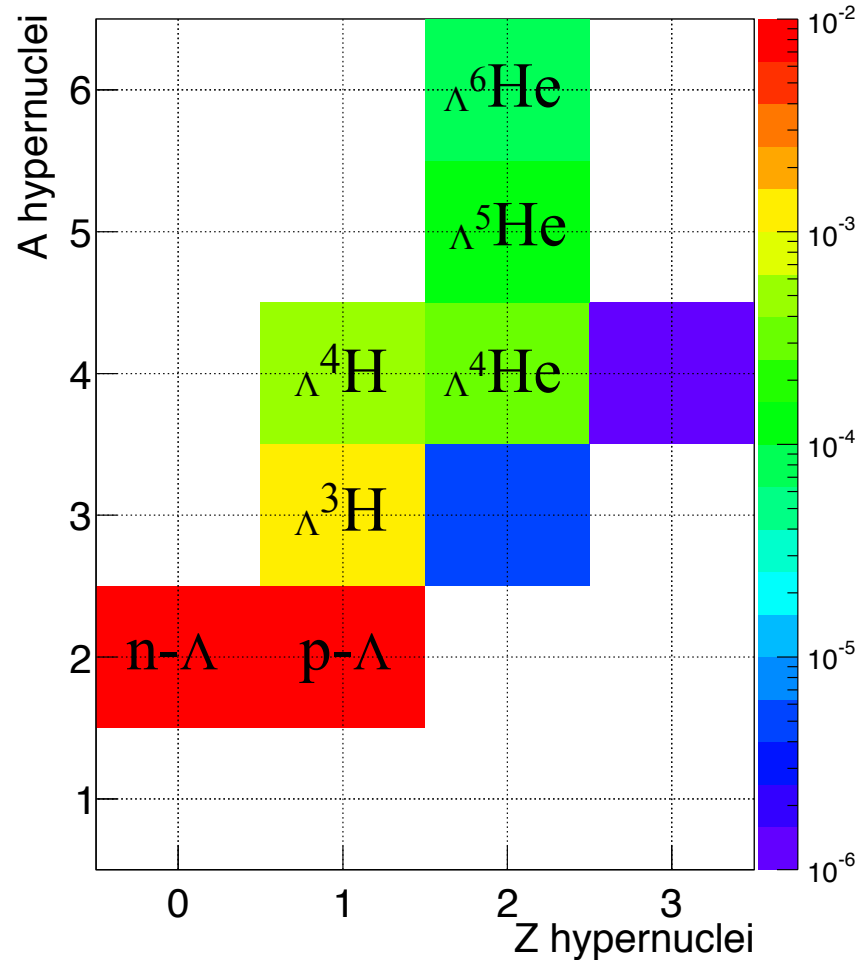
at

1.91 A.GeV

( $b < 6$  fm) -

$t_{\text{cluster.}} = 20$  fm/c

IQMD+SACA  $^{58}\text{Ni}+^{58}\text{Ni}$  at 1.93 A.GeV ( $b < 6$  fm,  $t_{\text{cluster.}} = 20$  fm/c) - soft+mdi, kaon pot.



Soft EOS  
with m.d.i.  
with Kaon pot.



# Another application of SACA : hypernuclei production

IQMD+SACA

$^{58}\text{Ni}+^{58}\text{Ni}$

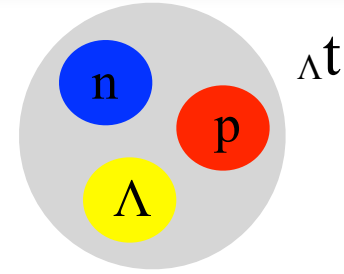
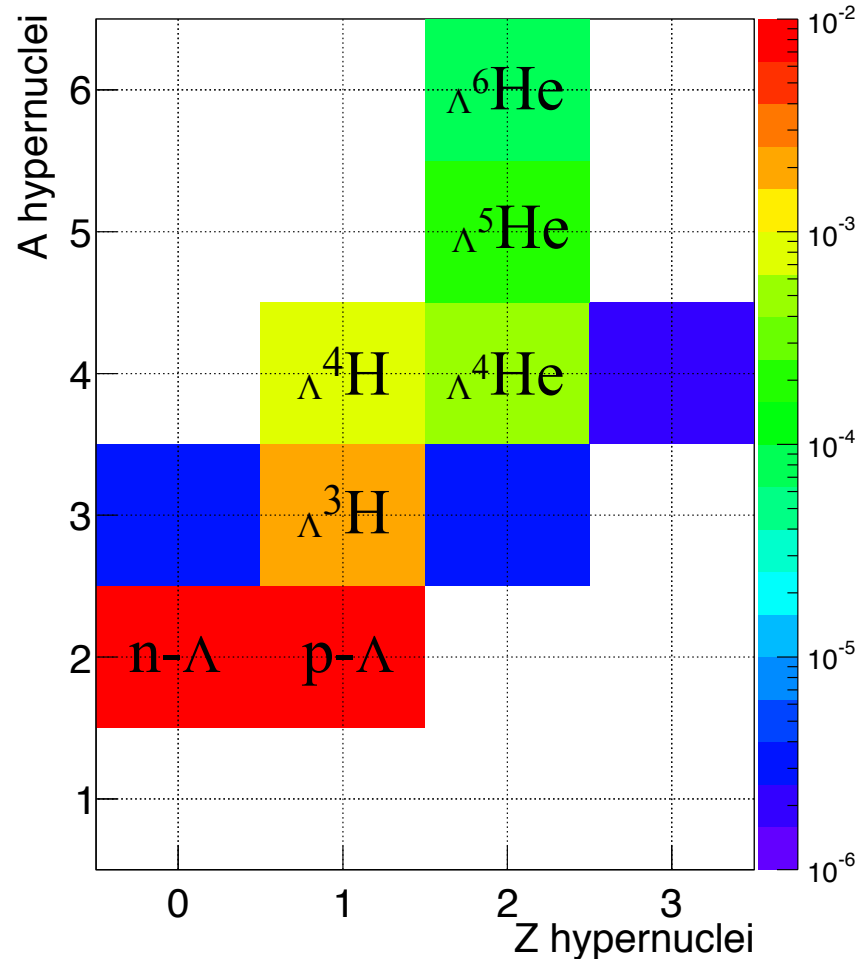
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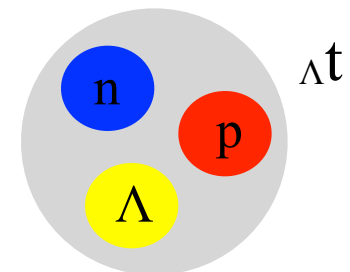
Soft EOS  
with m.d.i.  
no Kaon pot.





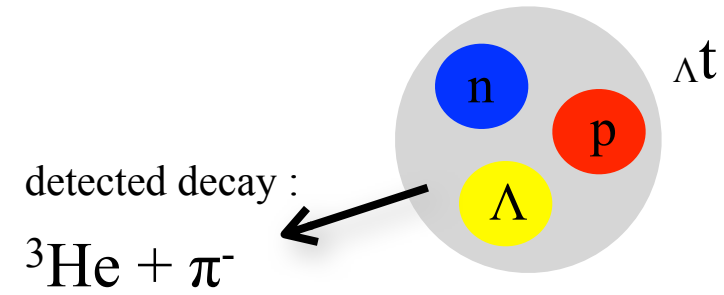


# Strong phase space constraints



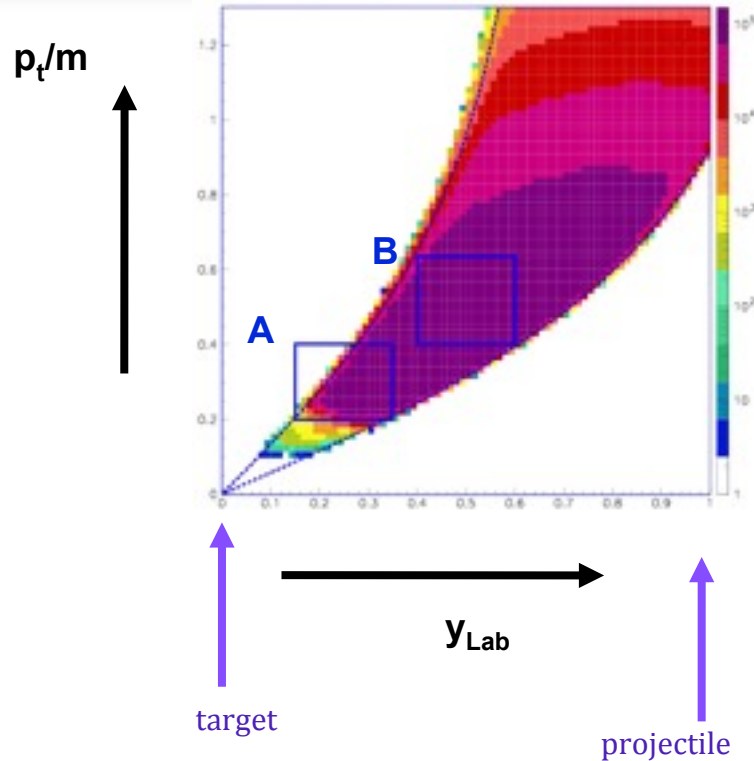


# Strong phase space constraints

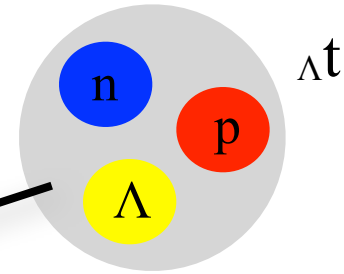
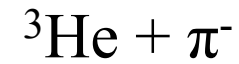


# Strong phase space constraints

FOPI Coll.  
Y. Zhang, Heidelberg



detected decay :



**Ni+Ni @ 1.91 A.GeV**

Preliminary

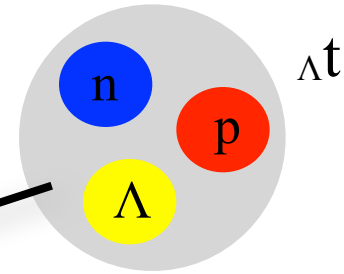
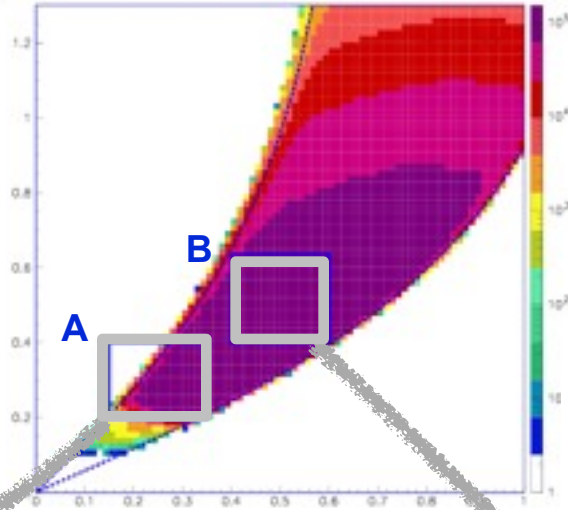


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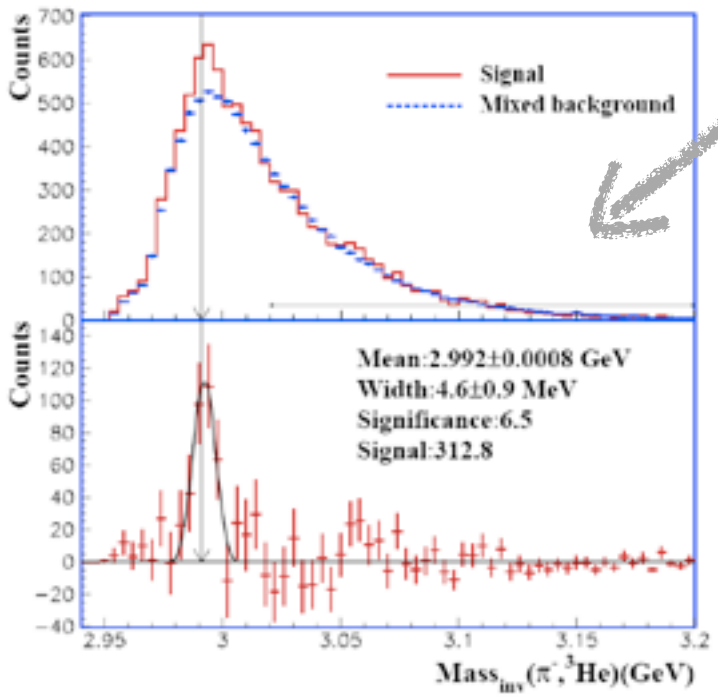
FOPI Coll.  
Y. Zhang, Heidelberg

Excess over combinatorial background only in region A

$p_t/m$



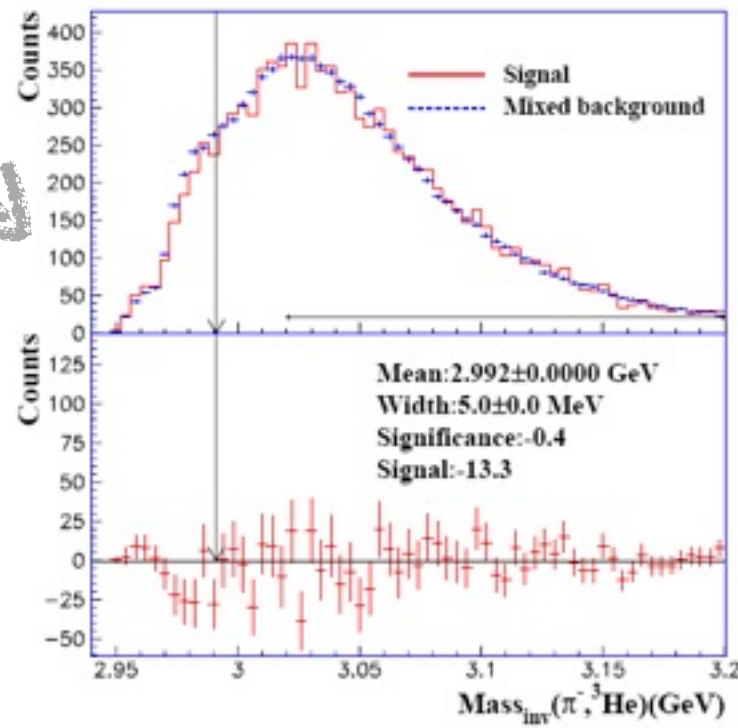
detected decay :  
 ${}^3\text{He} + \pi^-$



$y_{\text{Lab}}$   
target  
projectile

Ni+Ni @ 1.91 A.GeV

Preliminary



# Strong phase space constraints

IQMD+SACA

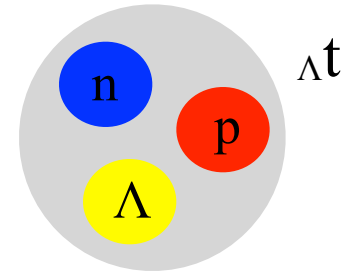
$^{58}\text{Ni}+^{58}\text{Ni}$

at

1.91 A.GeV

( $b < 6$  fm) -

$t_{\text{cluster.}}=20$  fm/c





# Strong phase space constraints

IQMD+SACA

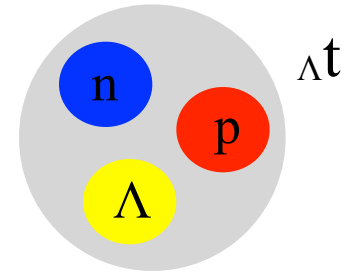
$^{58}\text{Ni}+^{58}\text{Ni}$

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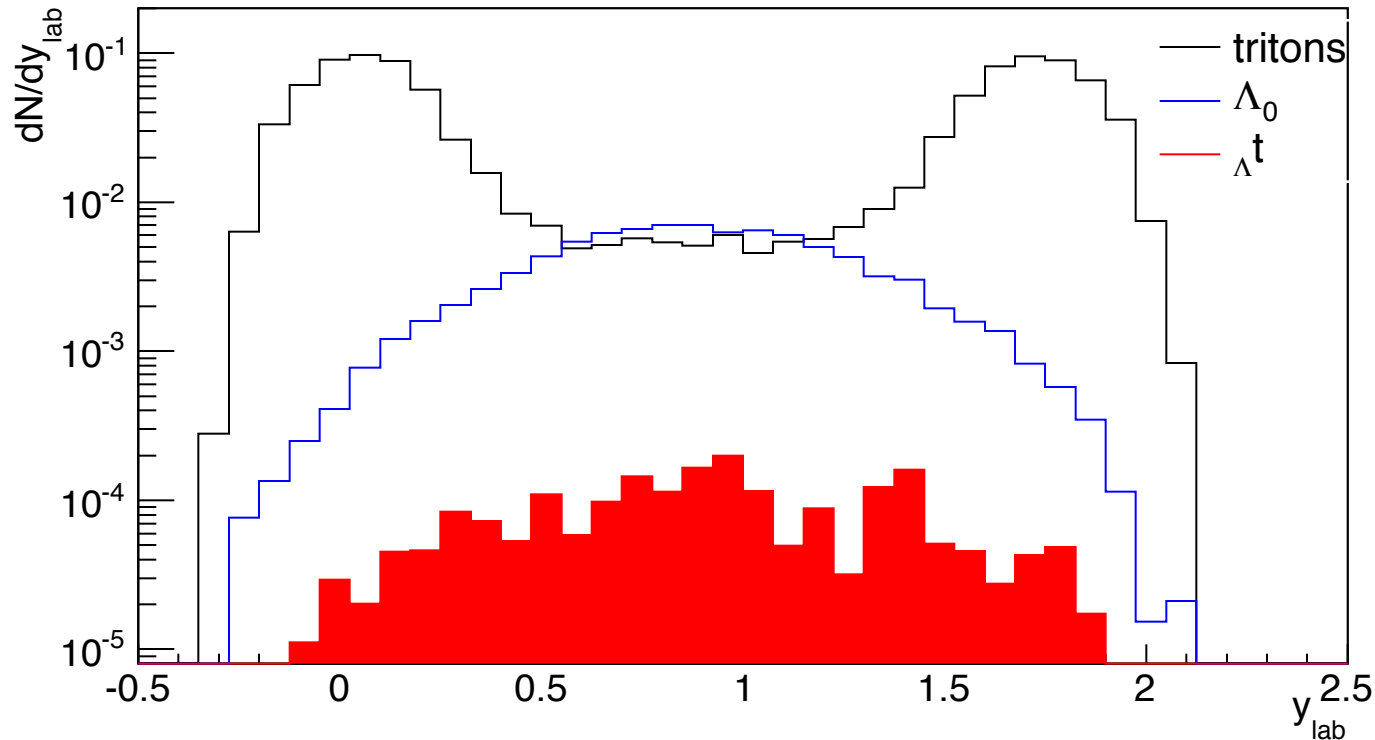
1.91 A.GeV

( $b < 6$  fm) -

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Soft EOS, no m.d.i., with Kaon pot.



# Strong phase space constraints

IQMD+SACA

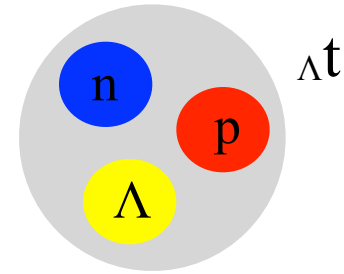
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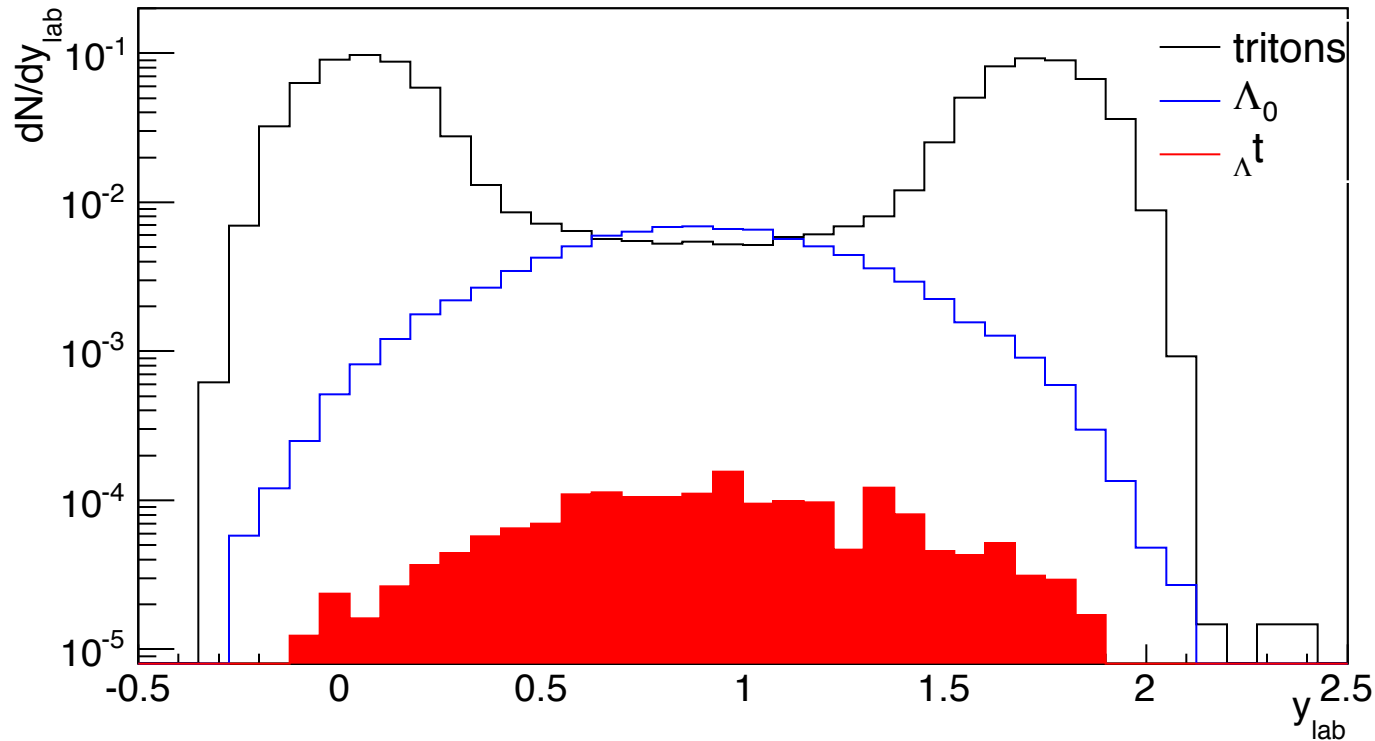
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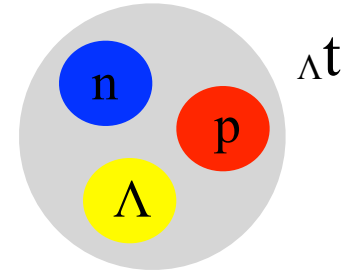
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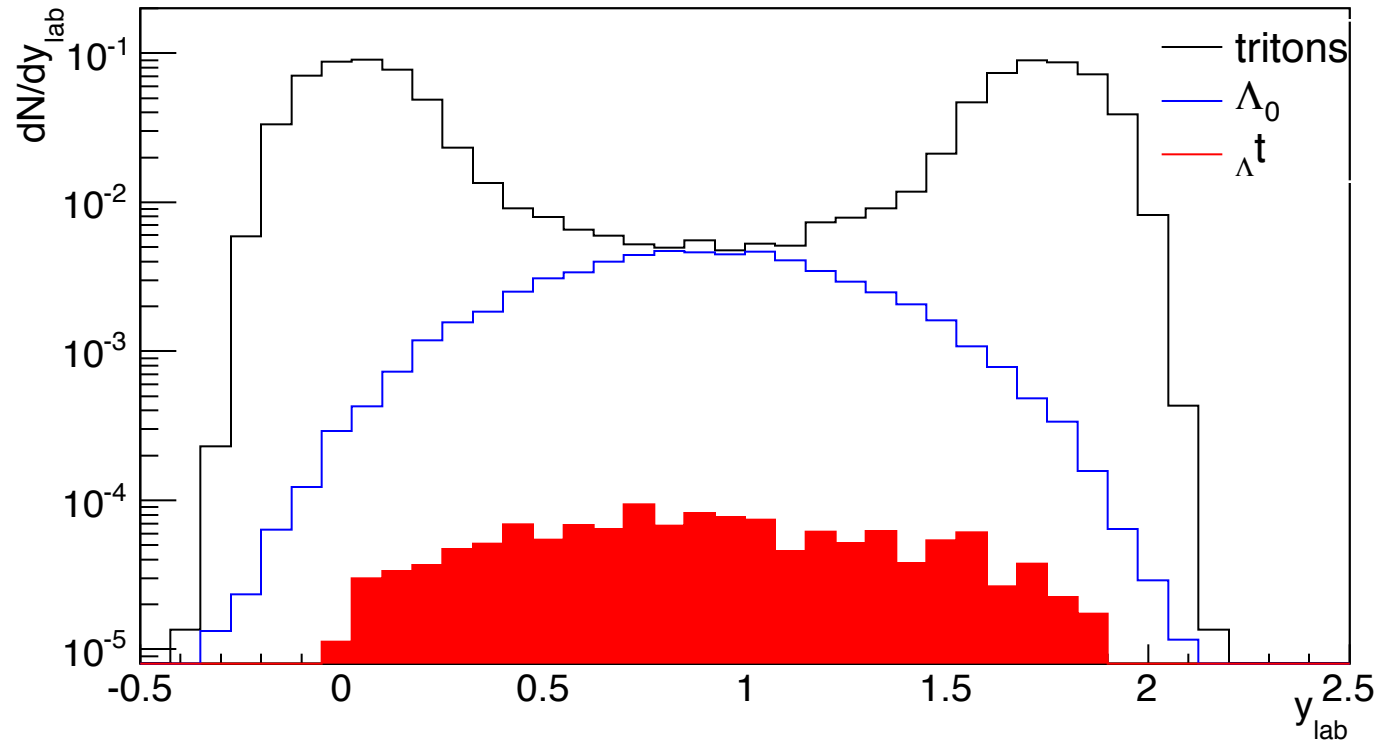
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Soft EOS, no m.d.i., with Kaon pot.







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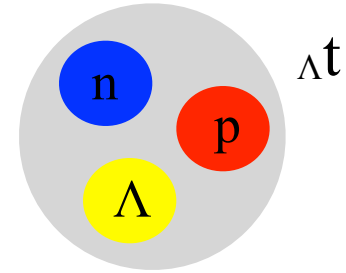
$^{58}\text{Ni}+^{58}\text{Ni}$

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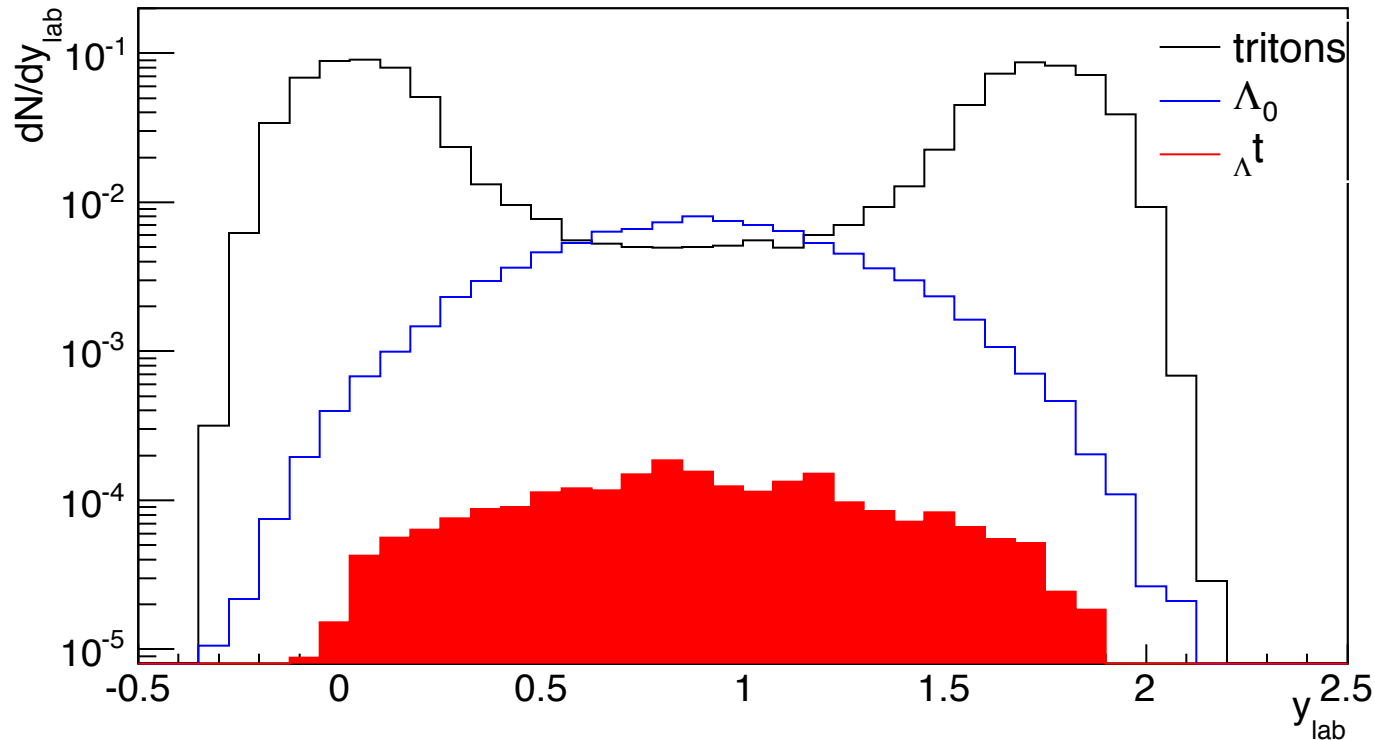
1.91 A.GeV

( $b < 6$  fm) -

$t_{\text{cluster.}}=20$  fm/c



Soft EOS, no m.d.i., no Kaon pot.

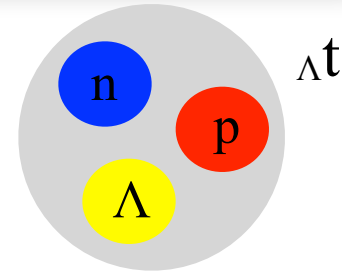


# Strong phase space constraints

IQMD+SACA

$^{58}\text{Ni}+^{58}\text{Ni}$  at 1.91 A.GeV

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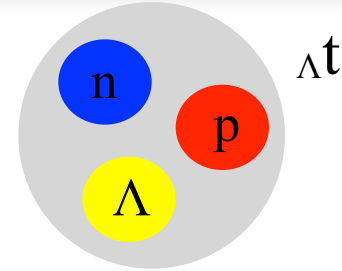
# Strong phase space constraints



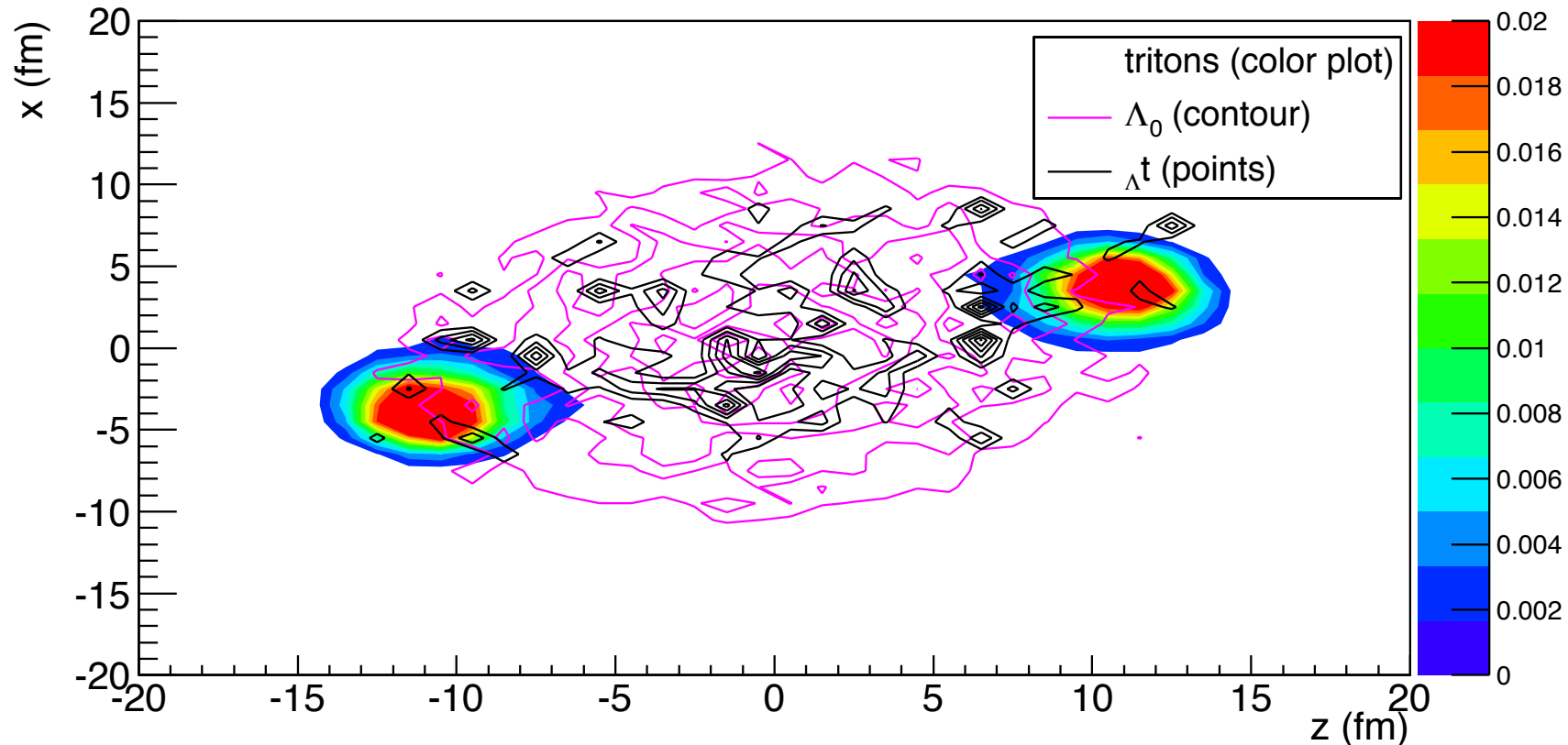
IQMD+SACA

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# Strong phase space constraints

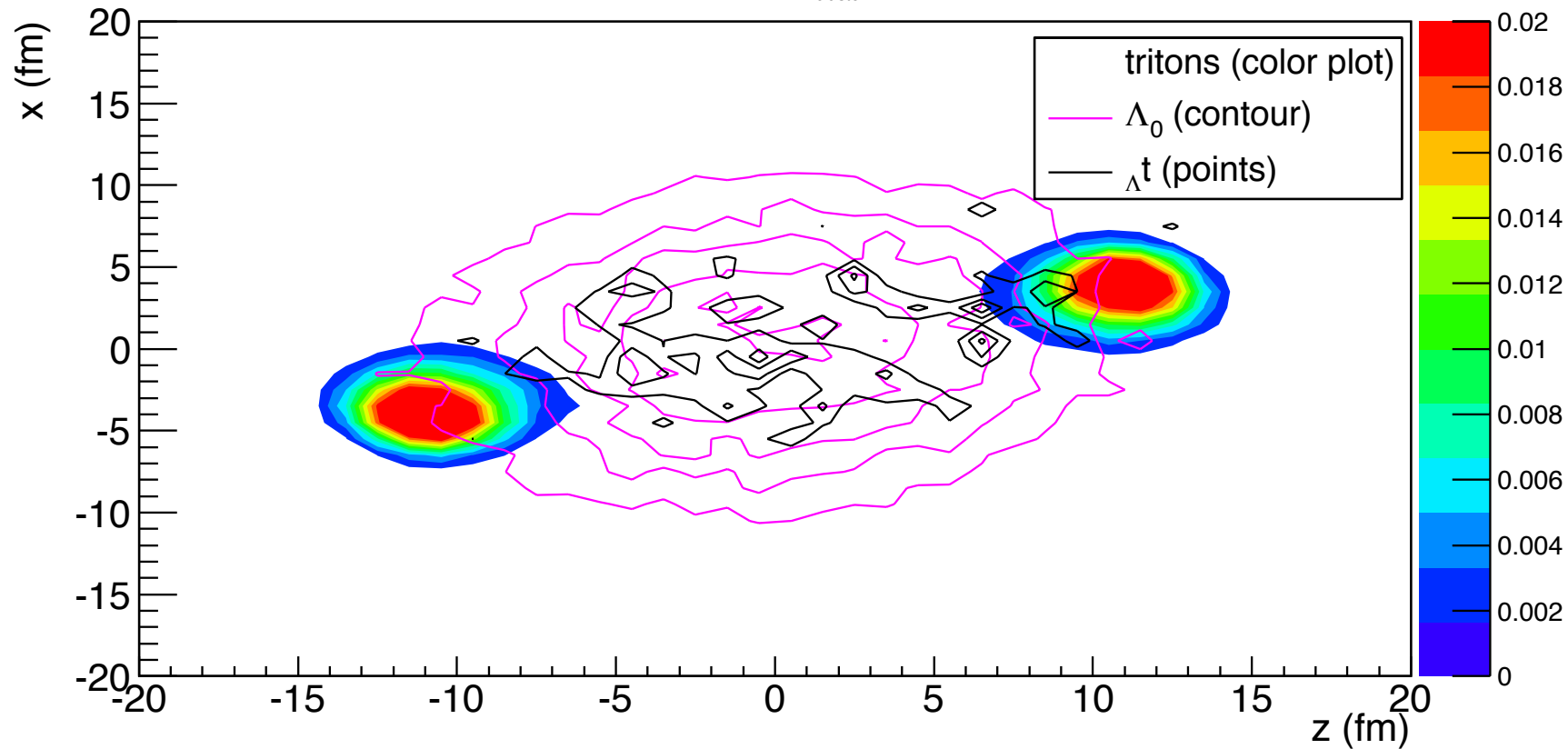
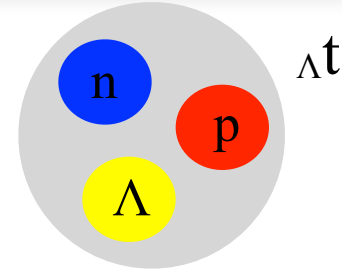


IQMD+SACA

$^{58}\text{Ni}+^{58}\text{Ni}$  at 1.91 A.GeV

( $b < 6$  fm) -  $t_{\text{cluster.}}=20$  fm/c

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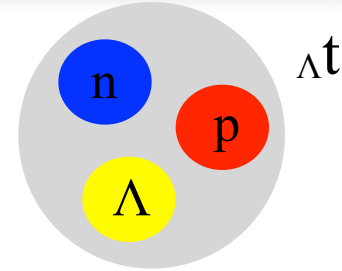
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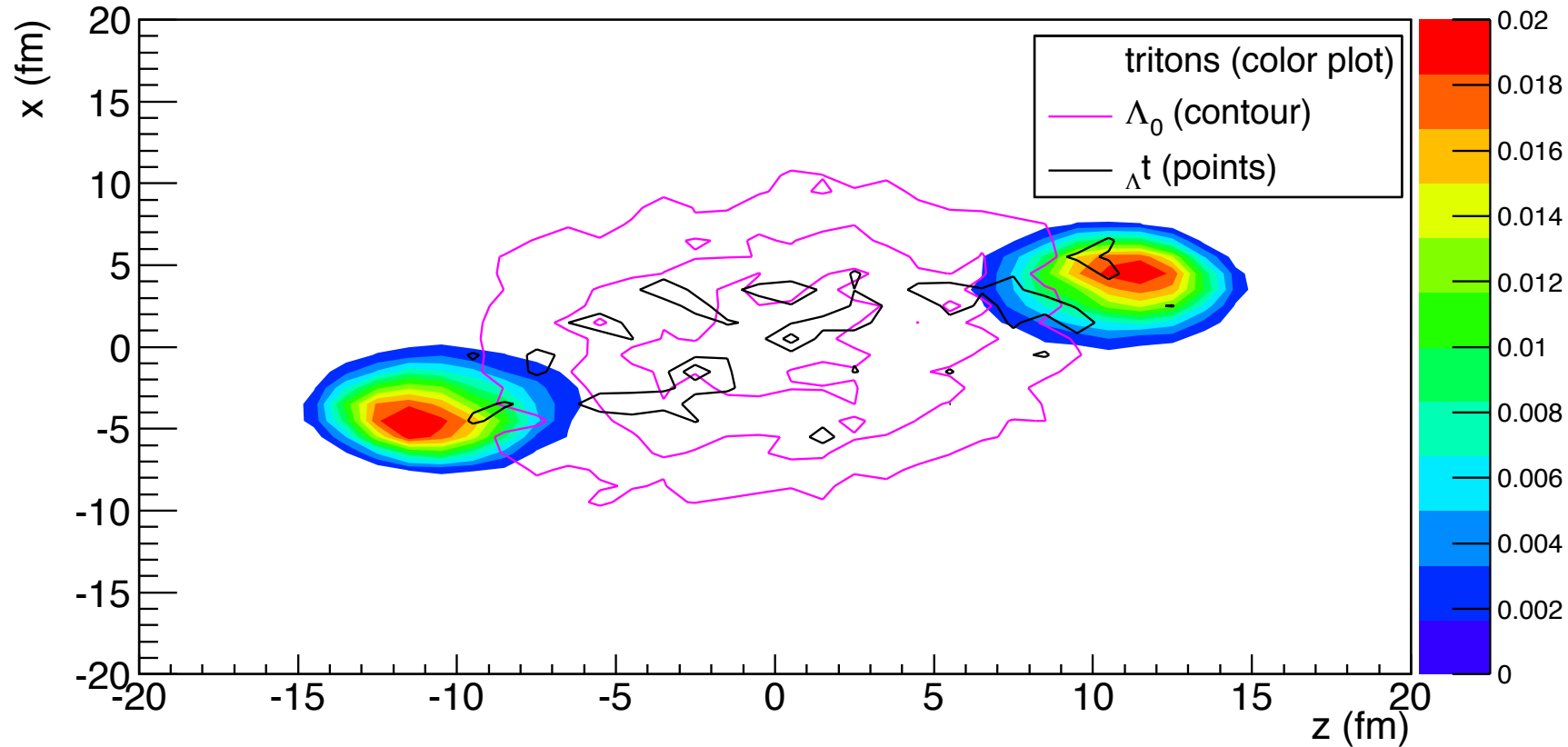
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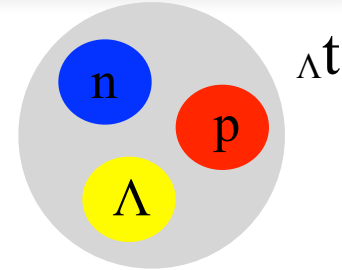
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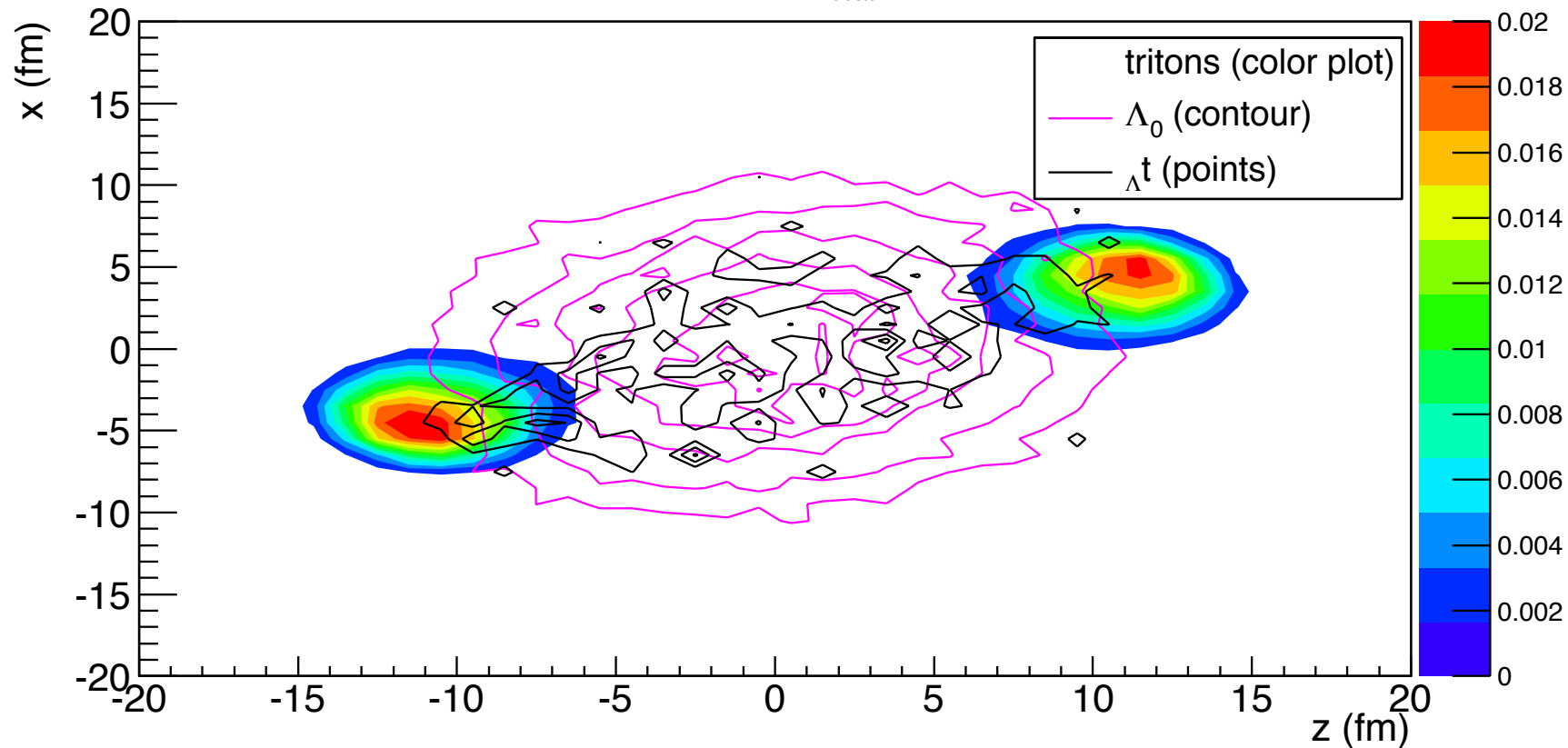
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- ❖ Supplying SACA with a more precise description of nuclei binding energy at abnormal density allows promising, realistic predictions of absolute isotope yields, and hypernuclei.
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- ❖ Within this model, isotope yields cannot inform on the high density dependence of the asymmetry energy.  $\Rightarrow$  better look at  $n / p$ ,  $K^+ / K^-$ ,  $\pi^+ / \pi^-$  yields/flows for that purpose.





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## Further developments:

After processing SACA, proceed:

- ❖ further decay of primary unstable isotopes like  $^8\text{Be}$ ,  $^5\text{He}$ , etc., which lifetime do not allow to detect them still bound in the detectors,
- ❖ allow early  $^3\text{He} + n \rightarrow ^4\text{He}$  according to its particularly high cross-section.
- ❖ secondary decay (evaporation code like GEMINI) of still excited clusters. Particularly relevant at intermediate energies ( $E_{\text{beam}}$  100 A.MeV down to the Fermi regime)
- ❖ for hypernuclei formation, refine  $\lambda$ -N potential in SACA or EOS/Kaon potential in IQMD in order to predict reasonably the measured cross-sections, and momentum distributions, which are very constraining.

# Density dependant pairing in SACA

Do the pairing and shell effect affect the primary fragments?

Probably **yes**, because:

✓ according to E. Khan et al., NPA 789 (2007) 94, pairing vanishes above  $T \approx 0.5\Delta_{\text{pairing}}$

$$\Delta_{\text{pairing}}(\rho_0) = 12 \text{ MeV}/\sqrt{A}; \Rightarrow \Delta B_{\text{pairing}}(^4\text{He})=12 \text{ MeV}, \Delta B_{\text{pairing}}(^3\text{He})=6.9 \text{ MeV}.$$

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E. Khan, M. Grasso and J. Margueron, in PRC 80 (2009) 044328 have derived the following function from the pairing potential  $V_{\text{pair}} = V_0 \{1 - \eta[\rho(r)/\rho_0]\} \delta(\mathbf{r}_1 - \mathbf{r}_2)$  within the Hartree-Fock-Bogolioubov method:

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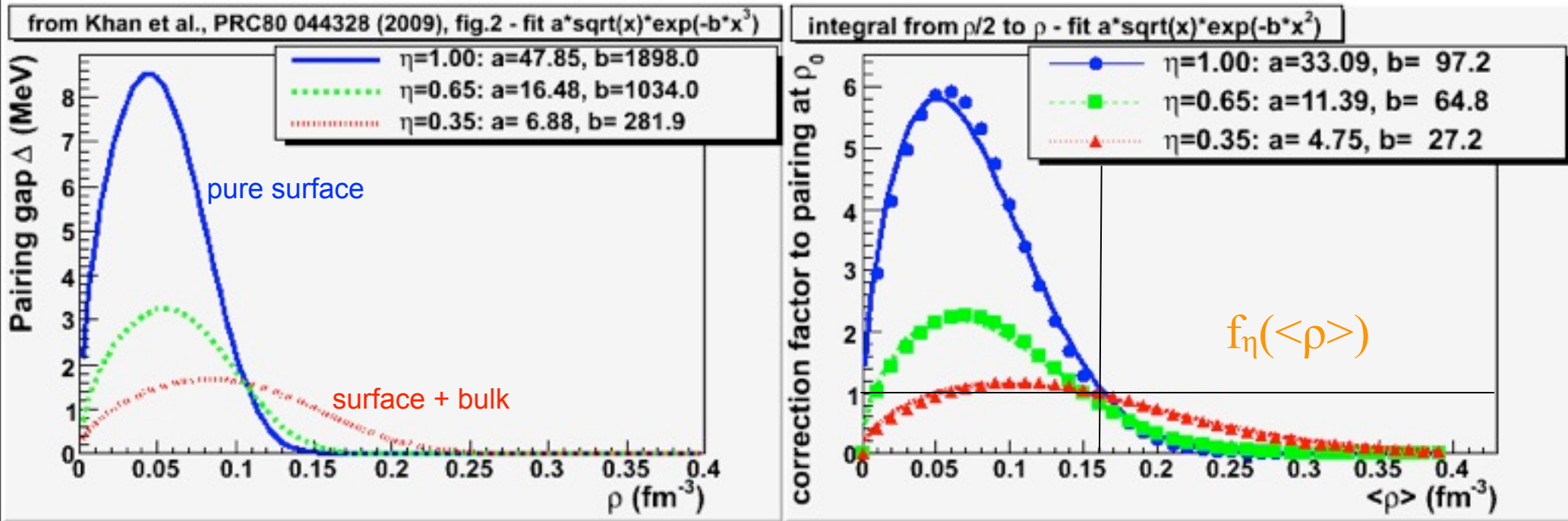
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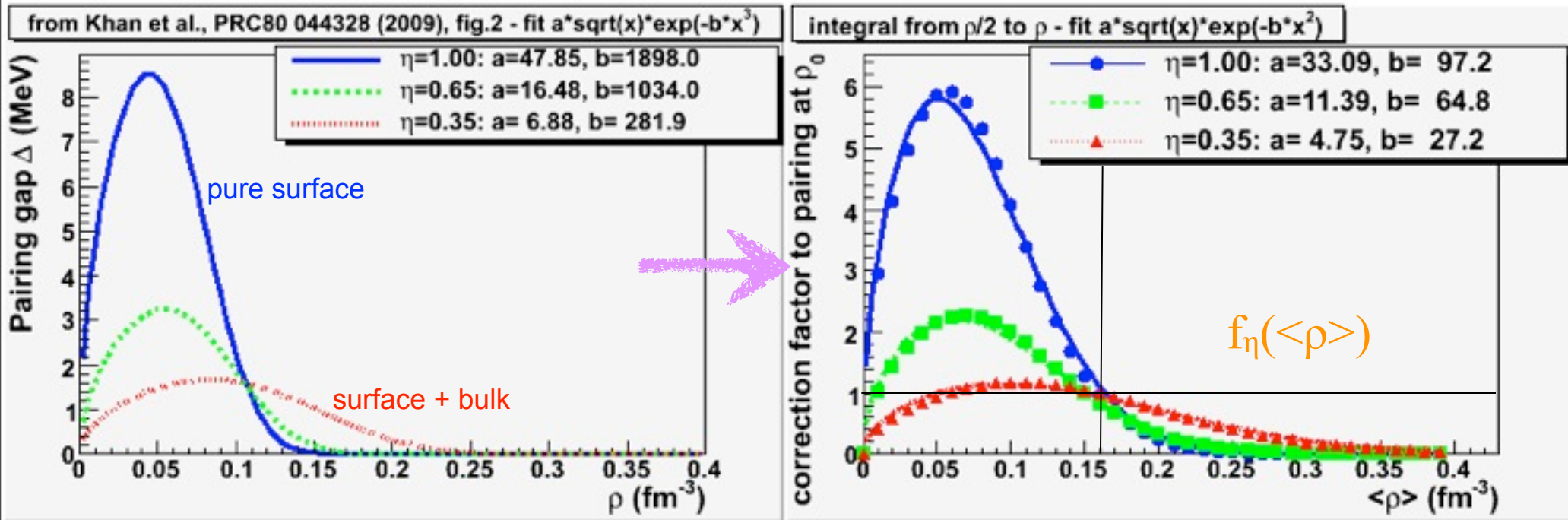
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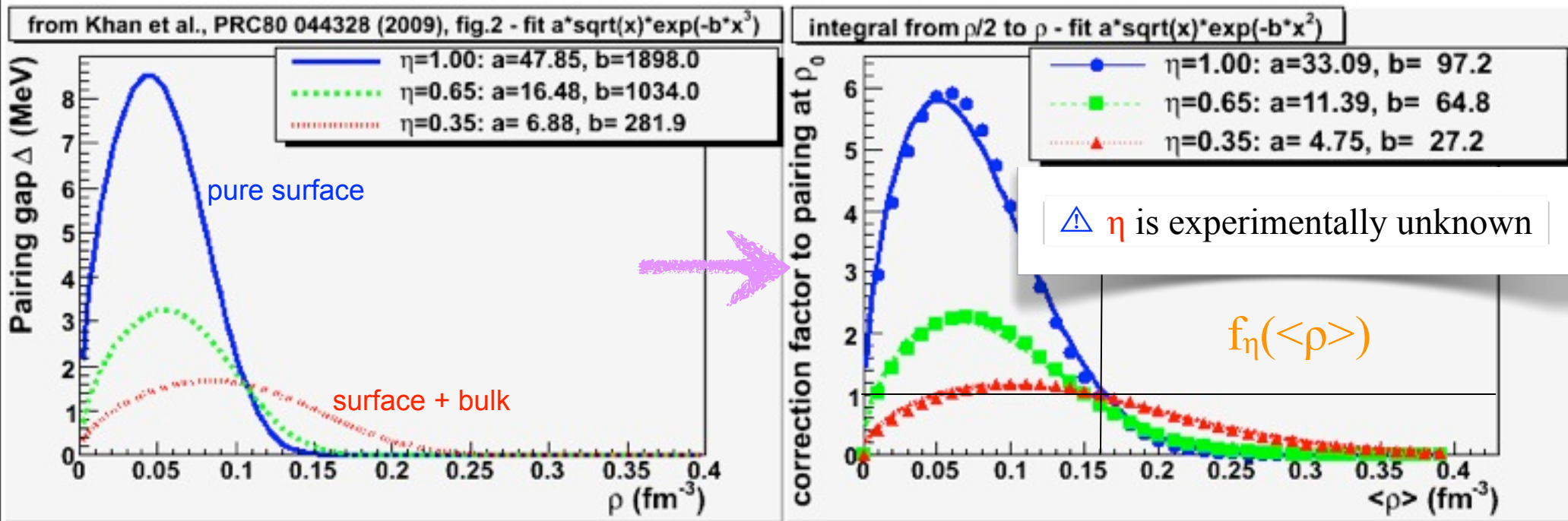
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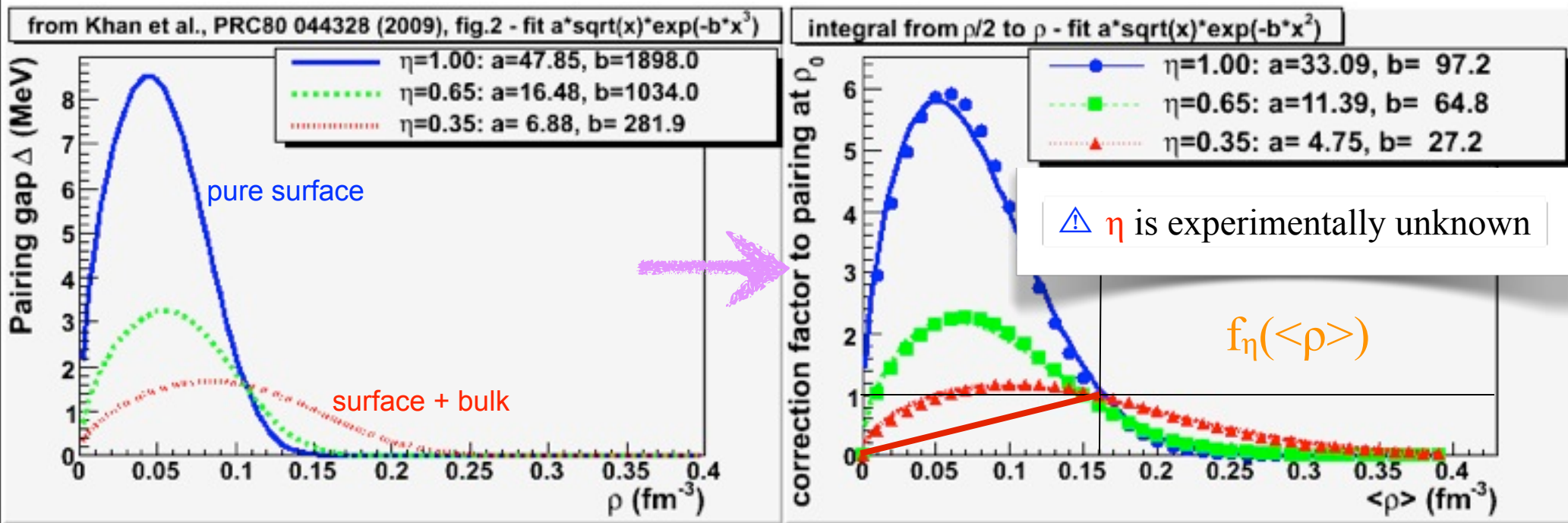
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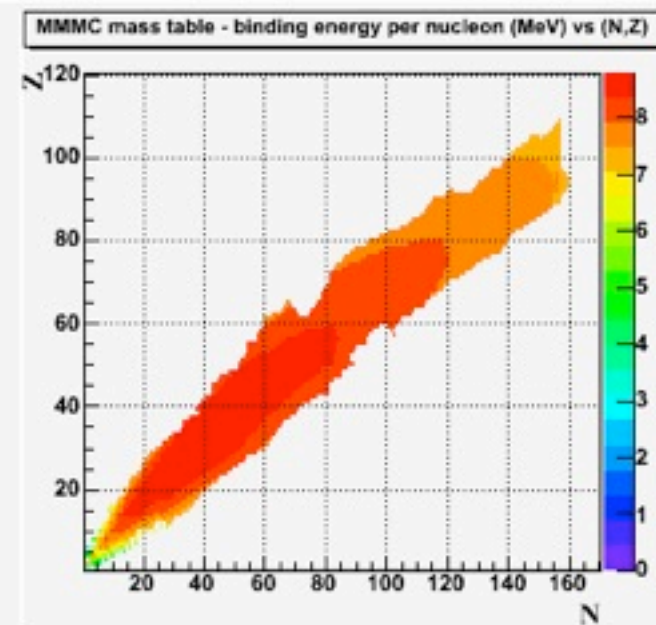
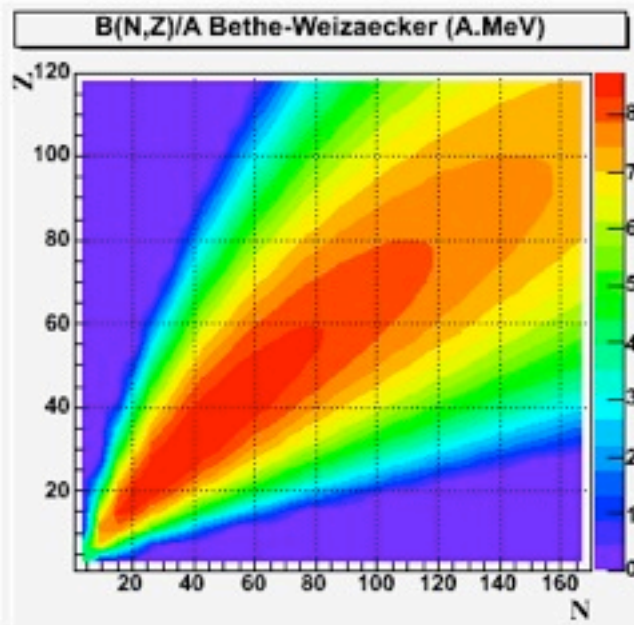
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☞ In order to account for all major structure effects which make the binding energy deviate from the liquid drop model, for each nucleus  $(N,Z)$ , what we call «pairing» binding energy will be the difference in binding energy between experimental measurements and the Bethe-Weizsäcker formula (without pairing).





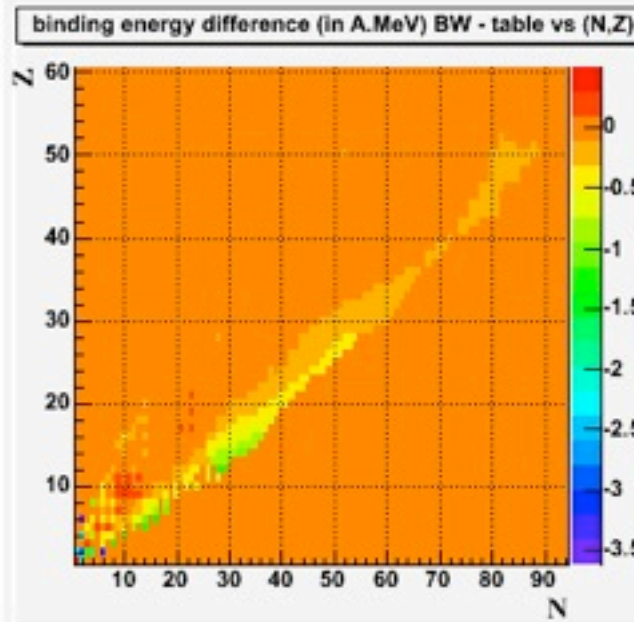
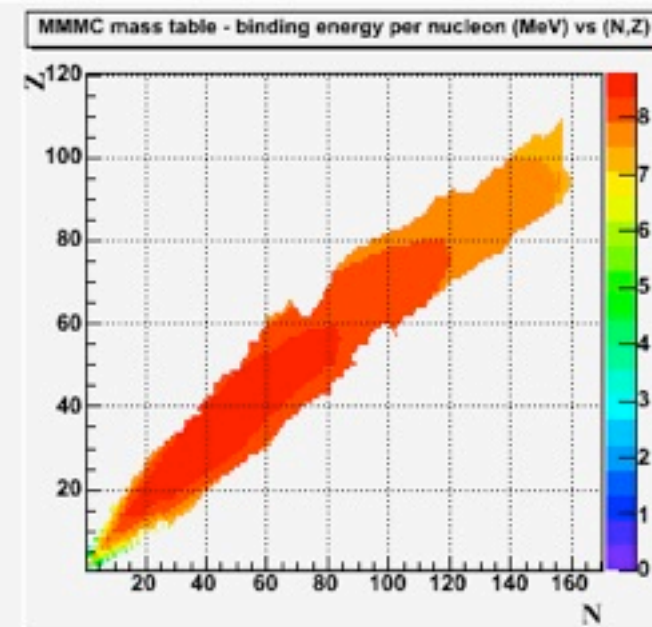
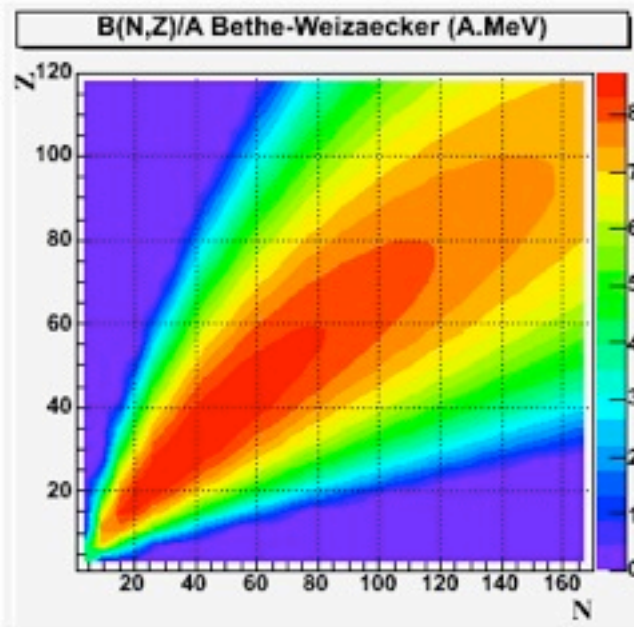
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→  $\Delta B_{\text{pairing}}(N,Z,\rho_0)$ .

And for a cluster at mean baryonic density  $\langle \rho_B \rangle$ , it will be

$\Delta B_{\text{pairing}}(N,Z,\rho_0) \times f_{\eta}(\langle \rho_B \rangle)$



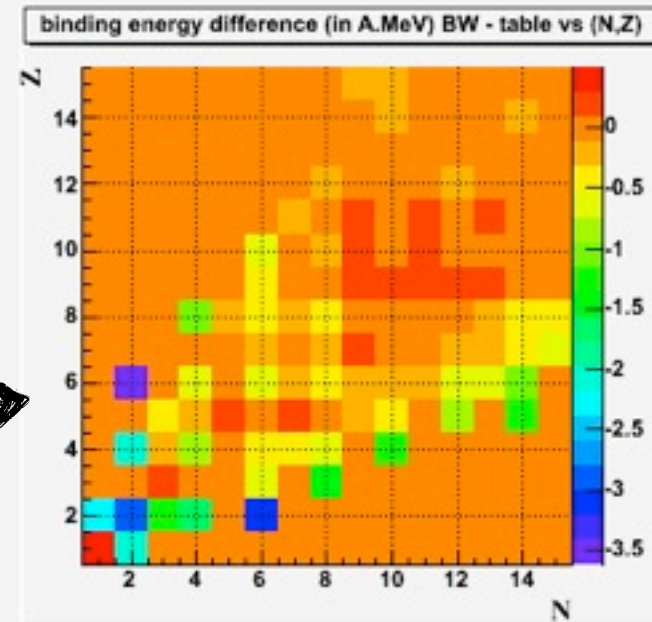
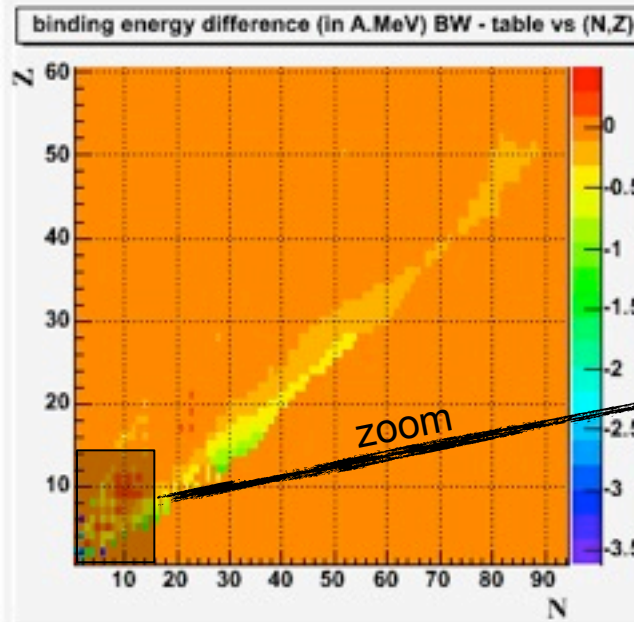
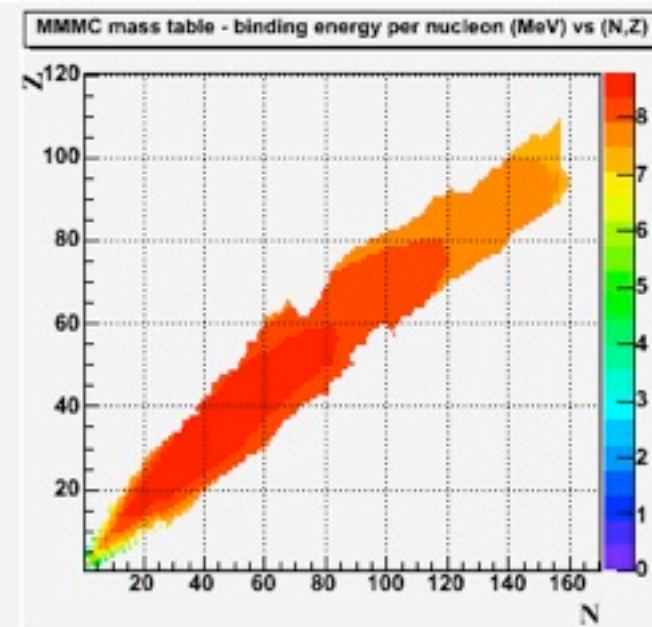
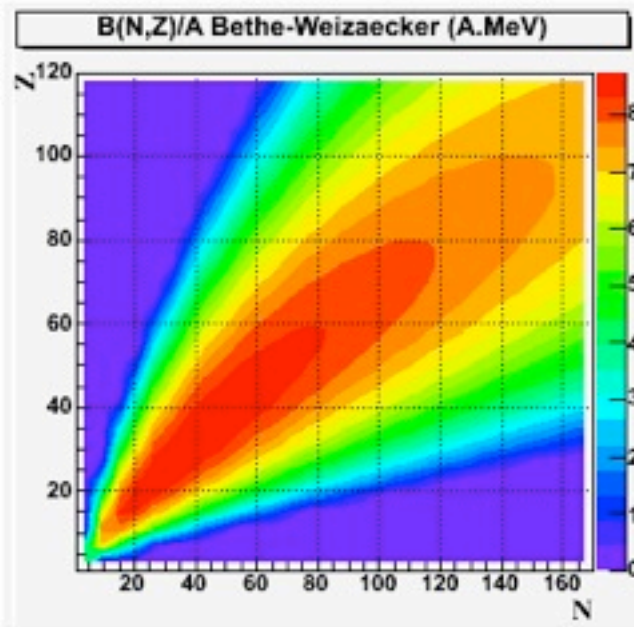
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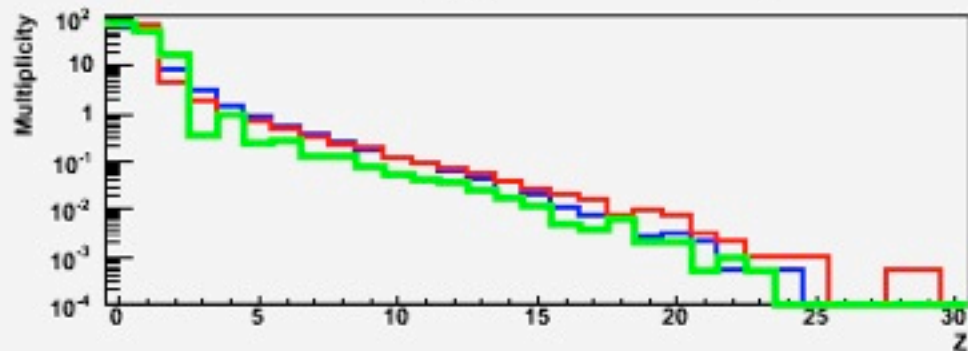


# SACA with asymmetry energy and pairing

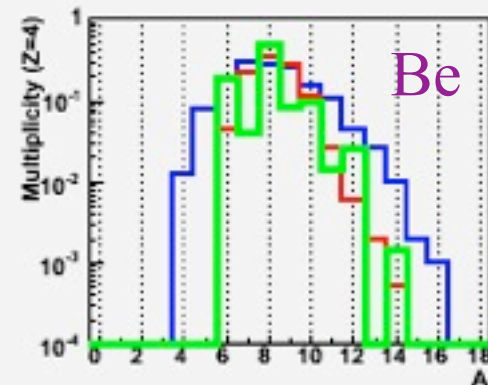
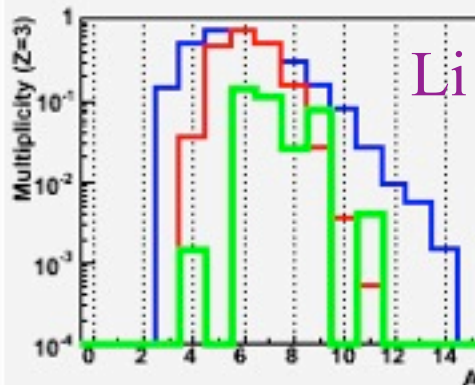
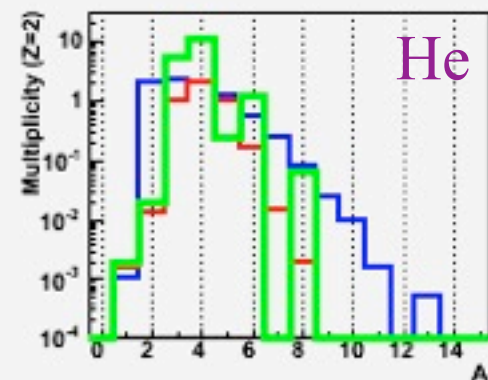
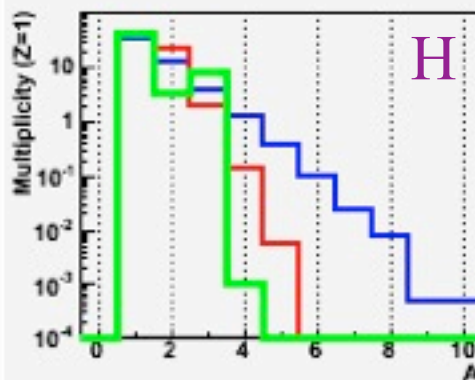
IQMD  $^{136}\text{Xe} + ^{112}\text{Sn}$  at 100 A.MeV,  $b=1$  fm,  $t_{\text{SACA}}=60$  fm/c

SACA version:

- $E_{\text{asy}}=0$ , no pairing
- $E_{\text{asy}}=32$  MeV ( $\gamma=1$ ), no pairing
- $E_{\text{asy}}=32$  MeV ( $\gamma=1$ ) +  $\eta_{\text{pairing}}=0.35$



$\eta = 0.35$   
surface & volume  
pairing

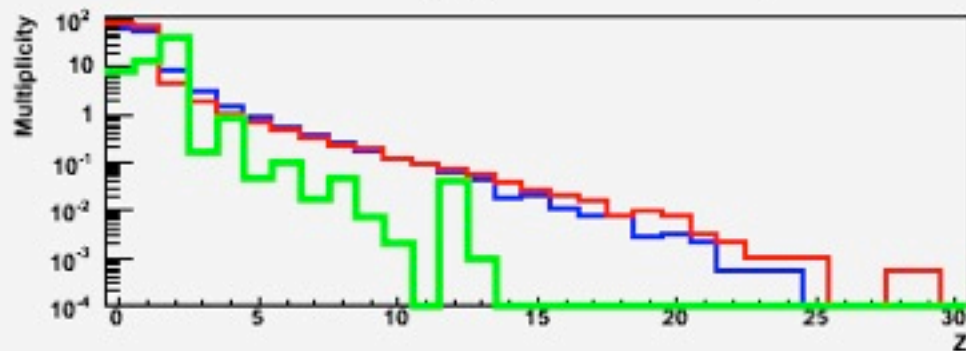


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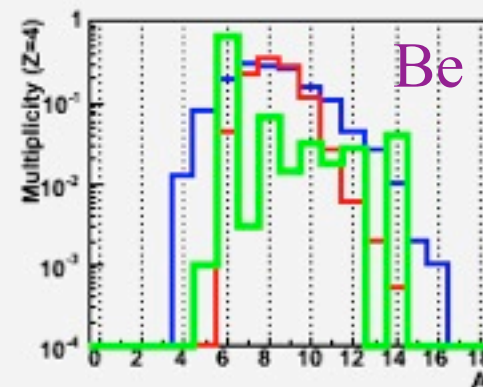
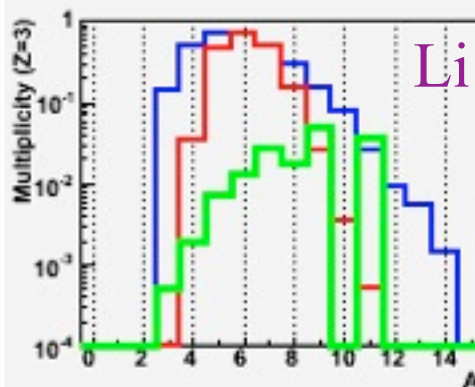
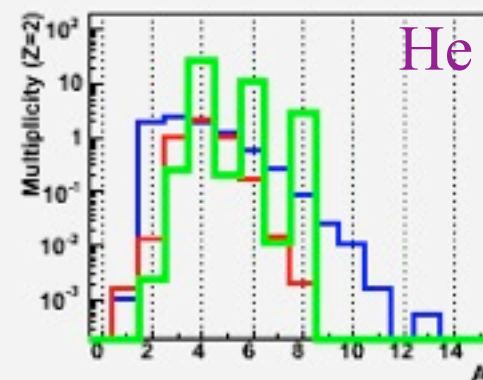
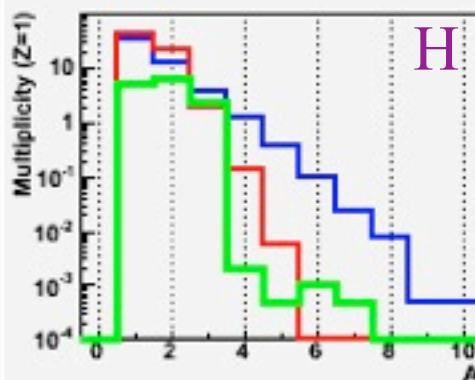
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$\eta = 1$   
pure surface  
pairing

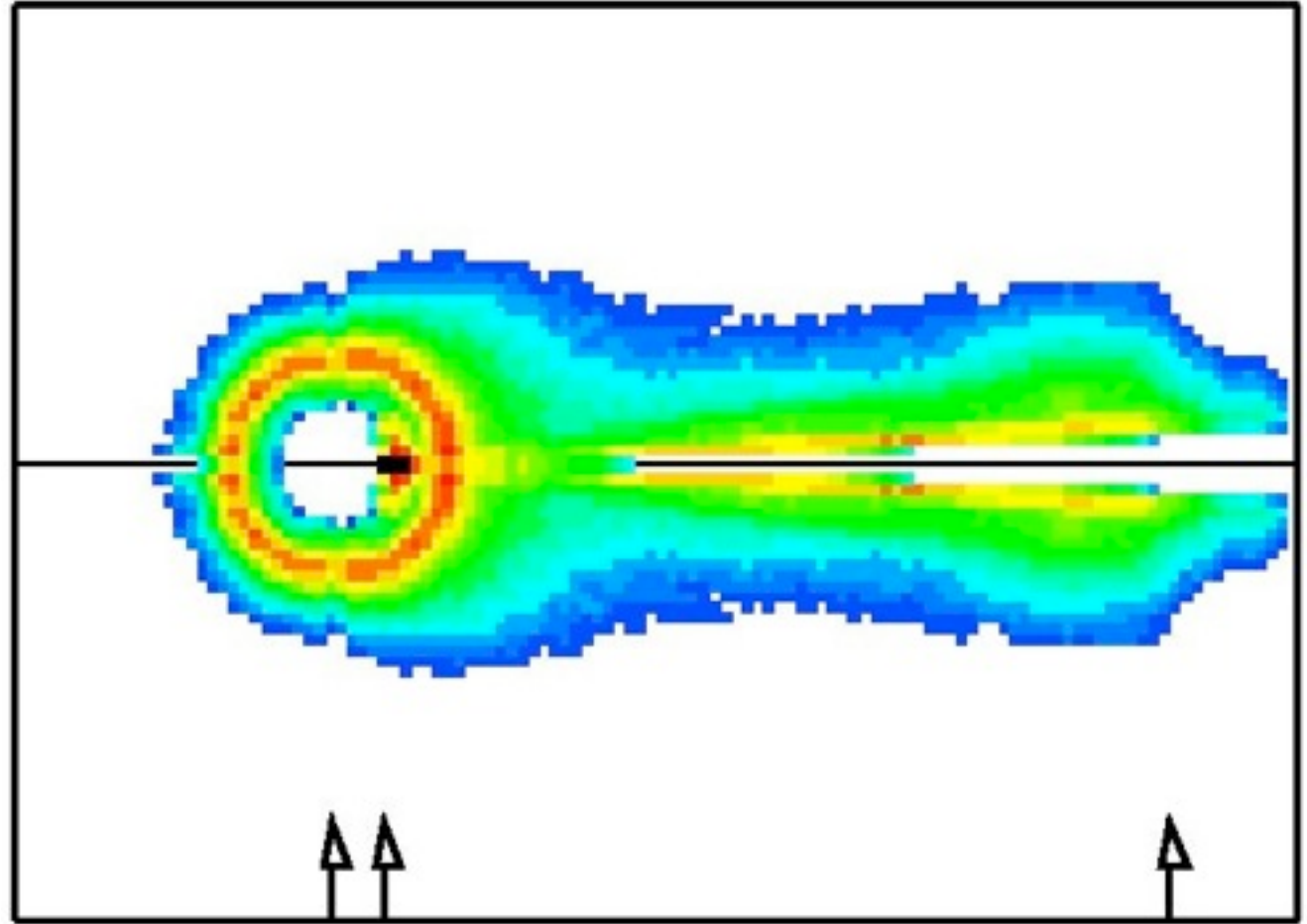




# Even in the spectator regime, we have to account correctly for isotope anisotropies

Illustration: alpha particles in C+Au at 300 A.MeV - peripheral collisions

$\gamma \beta_{\perp}$



$y = \text{th}^{-1} \beta_{\parallel}$

INDRA@GSI

J. Lukasik  
ALaDiN-INDRA Coll.

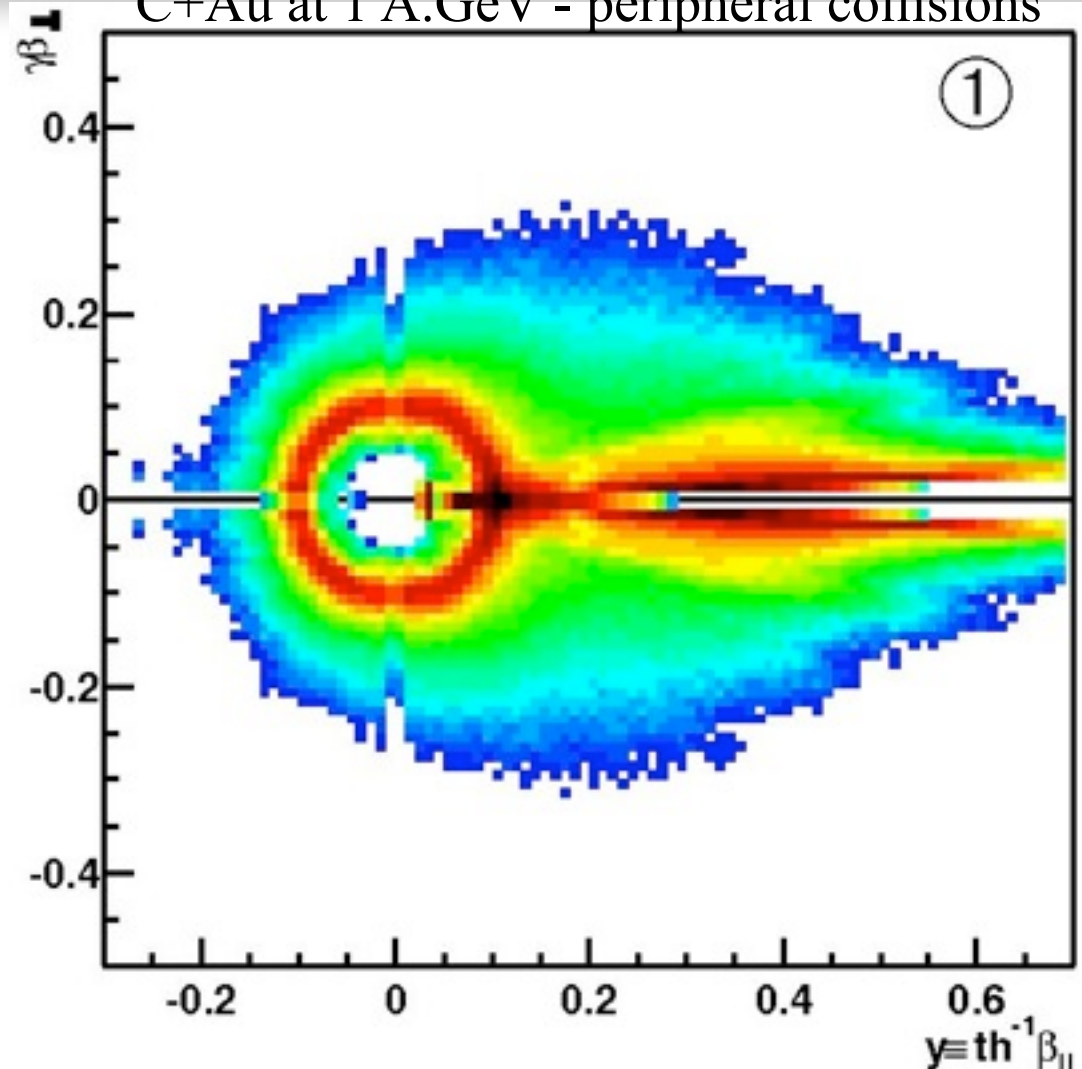




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Illustration: alpha particles in

C+Au at 1 A.GeV - peripheral collisions



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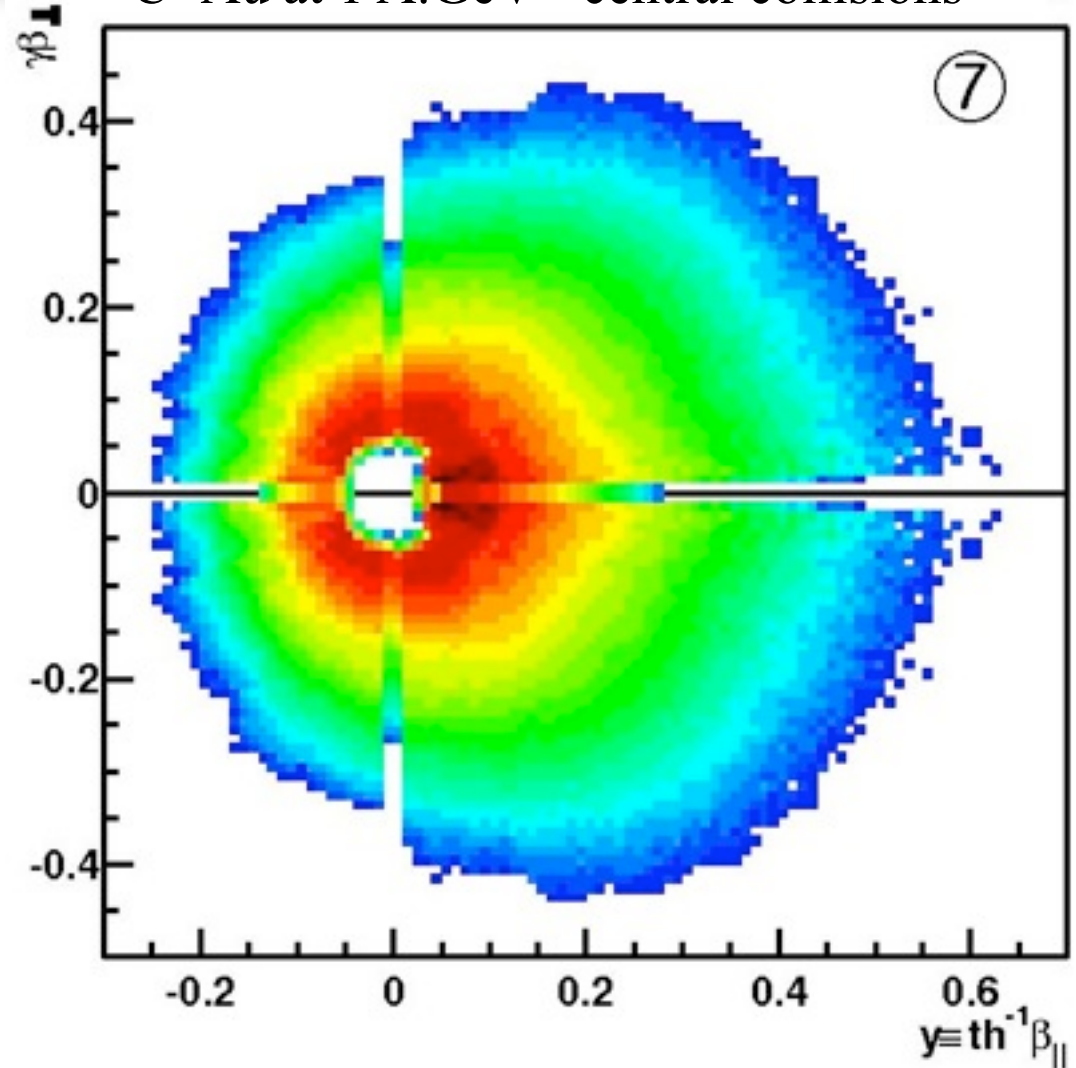




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Illustration: alpha particles in

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